

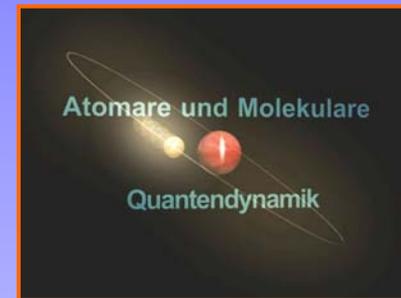
The World of Quantum Matter



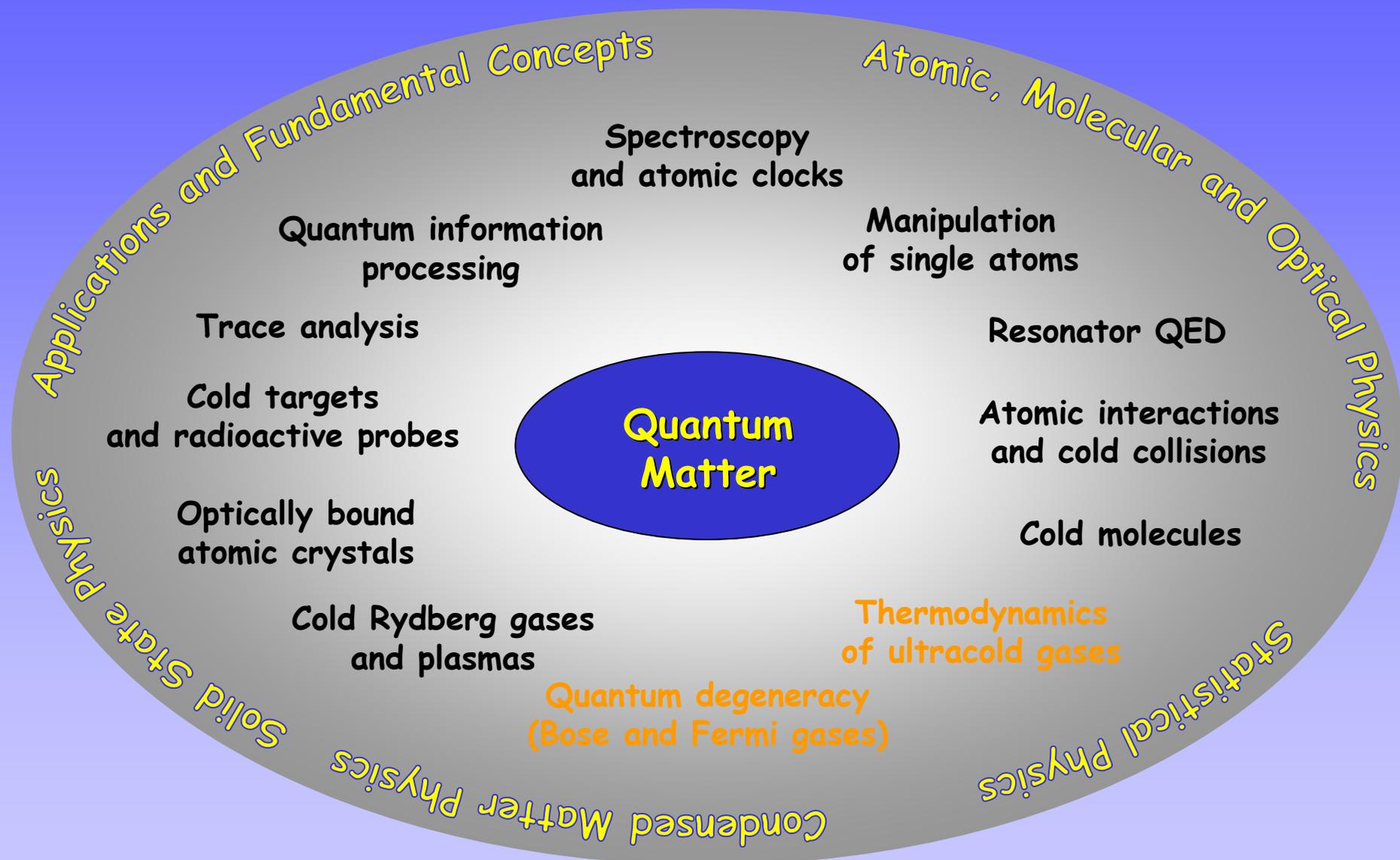
ALBERT-LUDWIGS-
UNIVERSITÄT FREIBURG

Atomic and Molecular Quantum Dynamics

Matthias Weidemüller
Albert-Ludwigs-Universität Freiburg



The World of Quantum Matter



Contents of the lectures

0. Primer on light-matter interactions
1. The way to absolute zero – cooling and trapping methods for atoms **Lecture 1**
2. Cold collisions
3. **Bose-Einstein condensation** **Lecture 2**
4. **Degenerate Fermi gases**
5. Cold Rydberg gases and plasmas **Lecture 3**
6. Ultracold molecules
7. Manipulation of single atoms **Lecture 4**
8. Cold atoms as targets for photon and particle beams

Bosons and Fermions

Quantum particles appear in two different “flavors” (quantum statistics) depending on their total angular momentum (spin):

integer spin

Bosons (e.g. photon)

half-integer spin

Fermions (e.g. electron, proton, neutron)

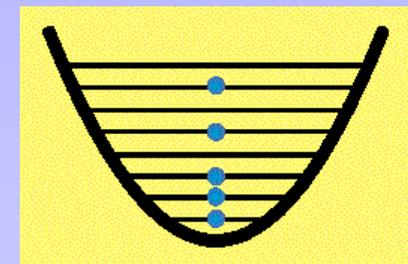
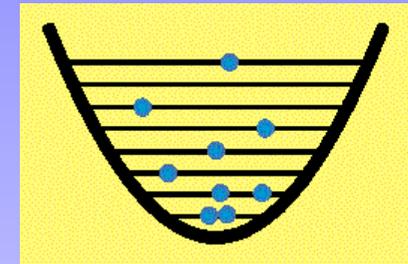
The spin determines the social quantum behaviour of these particles:

Bosons:

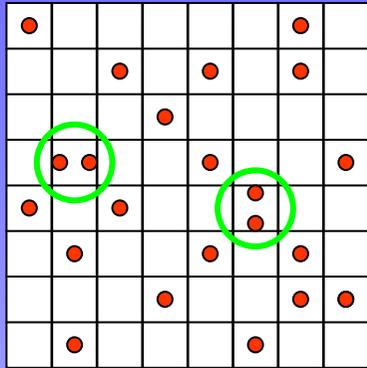
A quantum state can be occupied by an arbitrary number of bosons. If the state is already occupied by N bosons, the probability for the next boson to occupy the same state is N times enhanced.

Fermions:

A quantum state can never be occupied by more than one fermion.



Quantum statistics



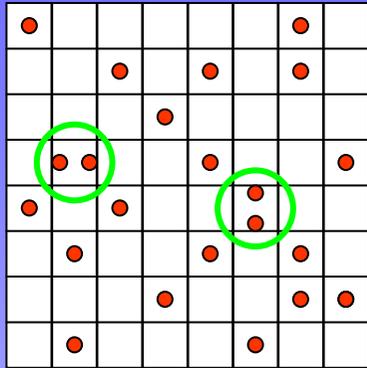
classical particles

$$p_i = \frac{1}{N}$$

occupation per state
follows simple Poissonian statistics

courtesy Rudi Grimm (Universität Innsbruck)

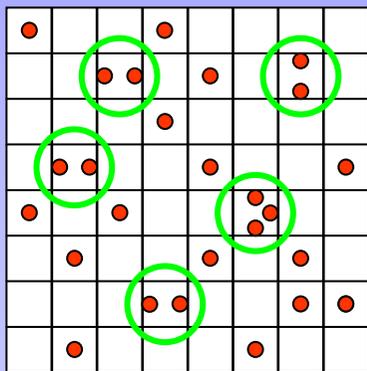
Quantum statistics



classical particles

$$p_i = \frac{1}{N}$$

occupation per state
follows simple Poissonian statistics



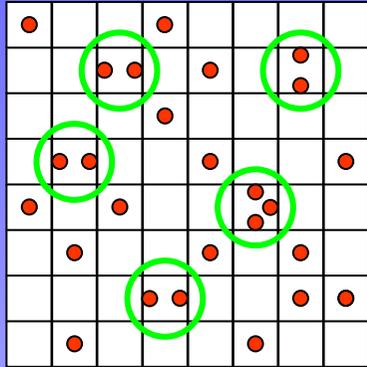
bosons

$$p_i \propto n_i + 1, \quad n_i = 0, 1, 2, \dots$$

bunching effect
(well known Hanbury-Brown-Twiss expt.)

courtesy Rudi Grimm (Universität Innsbruck)

Quantum statistics

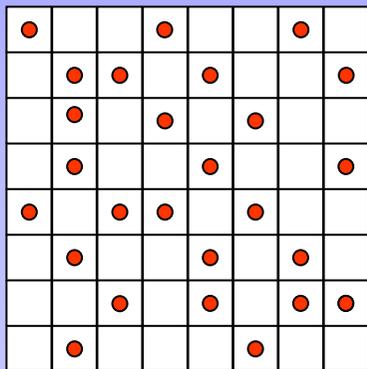


bosons

$$p_i \propto n_i + 1, \quad n_i = 0, 1, 2, \dots$$

bunching effect

(well known Hanbury-Brown-Twiss expt.)



fermions

$$p_i \propto n_i - 1, \quad n_i = 0, 1$$

Pauli's exclusion principle

courtesy Rudi Grimm (Universität Innsbruck)

Boltzmann factor



quantum states with different energies

$$p \propto e^{-E/k_B T}$$

Boltzmann factor

$$\text{Boltzmann's constant} \\ k_B = 1.3805 \times 10^{-23} \text{ J/K}$$



Ludwig Boltzmann

courtesy Rudi Grimm (Universität Innsbruck)

Distribution functions

Bose-Einstein statistics

$$f_{\text{BE}} = \frac{1}{e^{(E-\mu)/k_{\text{B}}T} - 1}$$

Boltzmann factor
together with
quantum statistics

$$f_{\text{cl}} = \frac{1}{e^{(E-\mu)/k_{\text{B}}T}}$$

classical limit

$$f_{\text{FD}} = \frac{1}{e^{(E-\mu)/k_{\text{B}}T} + 1}$$

Fermi-Dirac statistics

courtesy Rudi Grimm (Universität Innsbruck)

Thermodynamics

High temperature regime ($k_B T \gg \Delta E_{\text{q.m.}}$):

Each quantum mechanical state is occupied with a probability $\ll 1$
 \Rightarrow no difference between bosons and fermions (**Boltzmann statistics**)

Low temperature regime ($k_B T \lesssim \Delta E_{\text{q.m.}}$):

Each quantum mechanical state is occupied with a probability $\gtrsim 1$
 \Rightarrow **Bose-Einstein statistics** and **Fermi-Dirac statistics**

Level spacing $\Delta E_{\text{q.m.}}$ for particles with an average spacing d :

$$\Delta E_{\text{q.m.}} \sim p^2 / 2m \sim h^2 / 2md^2$$

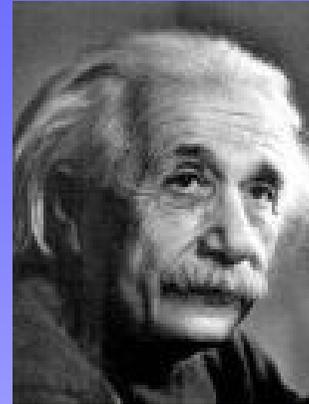
$$k_B T \lesssim \Delta E_{\text{q.m.}} \Leftrightarrow \boxed{n \Lambda_{\text{dB}}^3 \gtrsim 1}$$

thermal deBroglie wavelength $\Lambda_{\text{dB}} = (2\pi\hbar^2 / mk_B T)^{1/2}$

Prediction in 1925



Satyendranath Bose
(1894 - 1974)

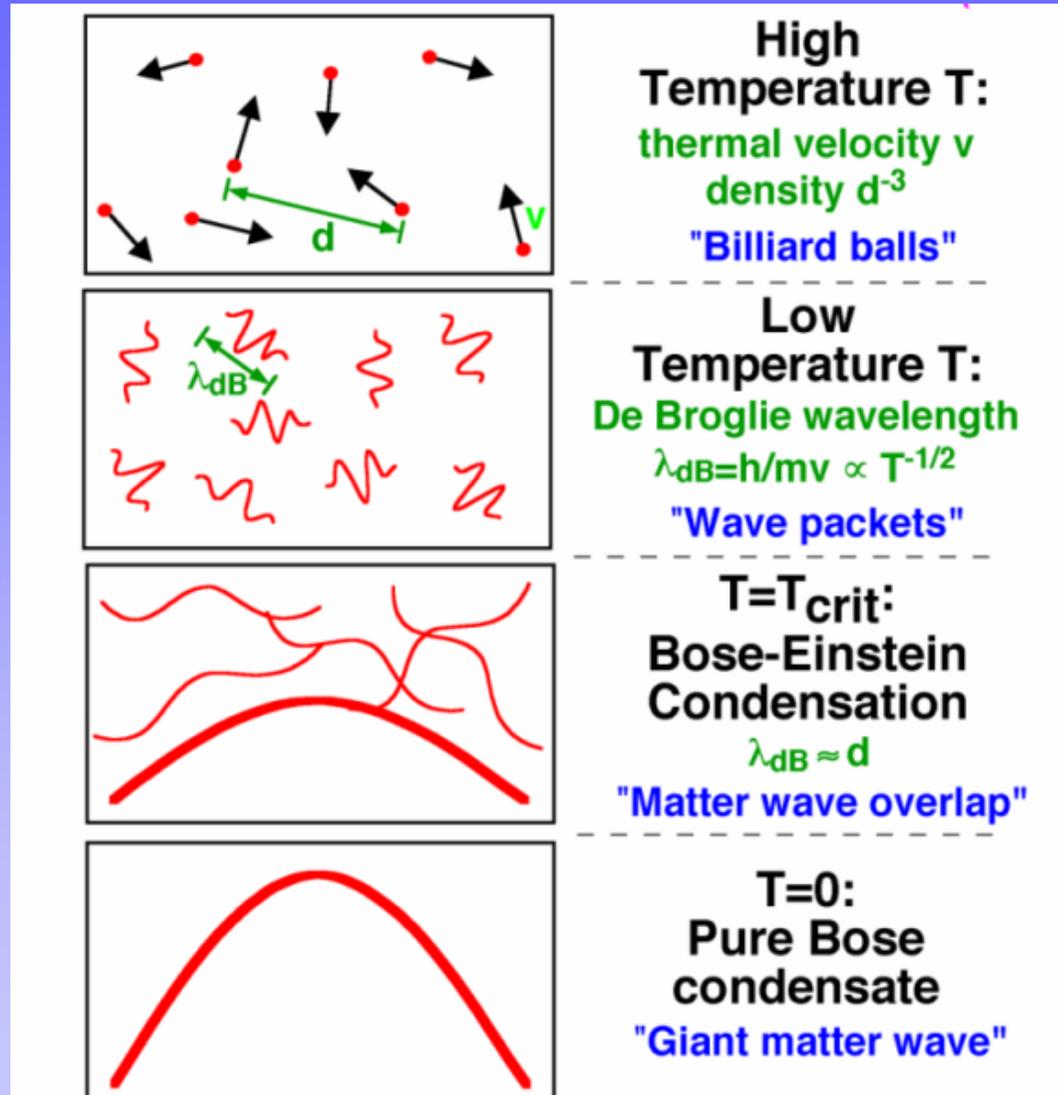


Albert Einstein
(1879 - 1955)

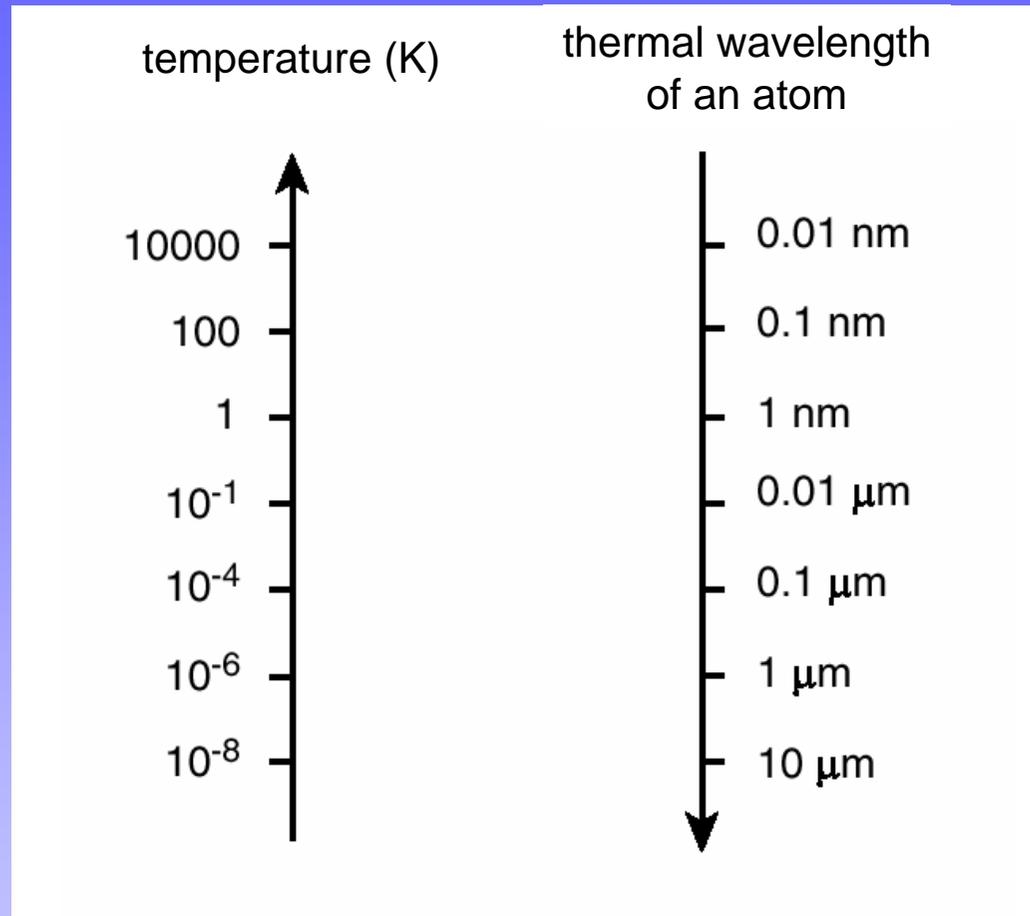
An **“ideal”** gas of **Bosons** shows a **phase transition** at sufficiently **low temperatures**.

The gas condenses into the lowest available quantum state and forms a **macroscopic quantum object**:
the **Bose-Einstein condensate**.

Bose-Einstein condensation (BEC)



Thermal deBroglie wavelength

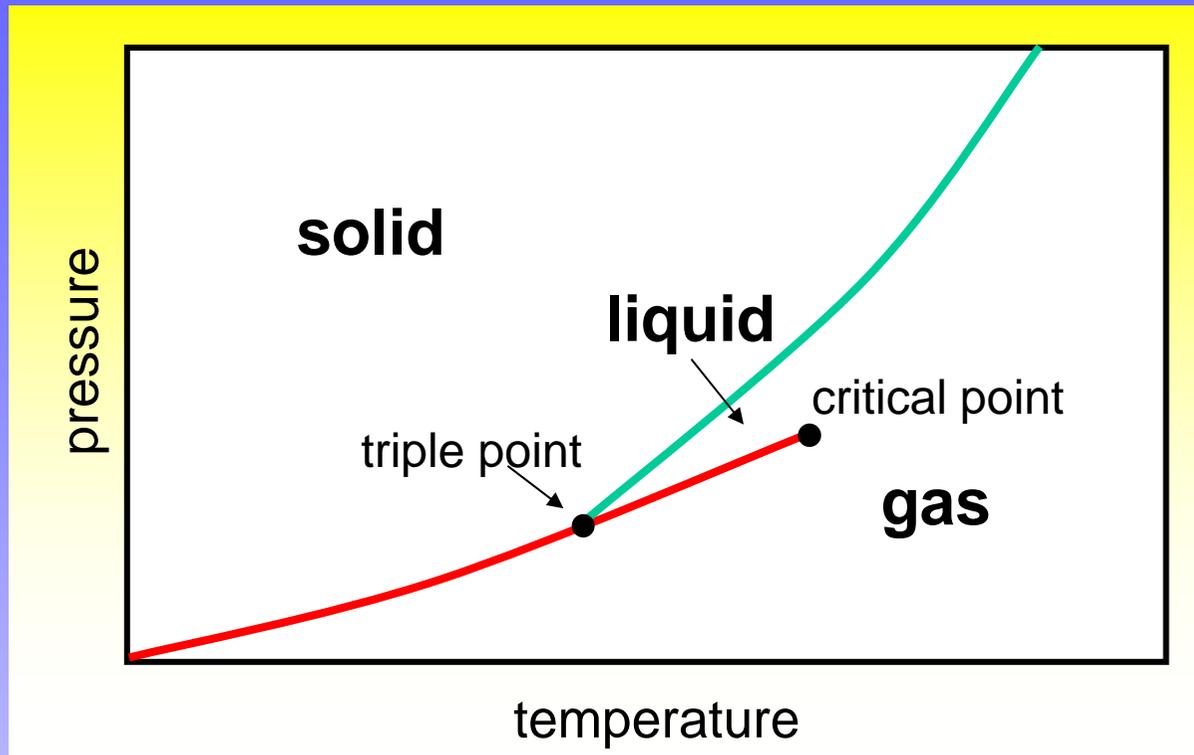


At room temperature (300 K):

Bose condensation would occur at a density larger than $(0.1 \text{ nm})^{-3} = 10^{24} \text{ atoms/cm}^3$
(typical density in a solid sample !)

Why we do not observe Bose condensation in real life ?

Phase diagram



At low temperatures, the thermal equilibrium state of every system is the **solid phase** (even down to $T = 0$)

Bose-Einstein condensation of a gas can only occur as a **metastable phase** at low densities ($\sim 10^{14}$ atoms/cm³) to prevent 3-body-recombination (equivalent to a saturated vapor)

⇒ **ultralow temperatures required** ($\sim 1 \mu\text{K}$)

Are atoms bosons?

Atoms and molecules are composite particles, composed of fermionic elementary particles (electrons, protons, neutrons).

Total spin is integer (total number of electrons, protons and neutron is even)

⇒ atoms and molecules are ***bosonic***

Total spin is half-integer (total number of electrons, protons and neutron is odd)

⇒ atoms and molecules are ***fermionic***

Are atoms bosons?

Atoms and molecules are composite particles, composed of fermionic elementary particles (electrons, protons, neutrons).

Total spin is integer (total number of electrons, protons and neutron is even)

⇒ atoms and molecules are ***bosonic***

Total spin is half-integer (total number of electrons, protons and neutron is odd)

⇒ atoms and molecules are ***fermionic***

Under which conditions is a composite particle a composite particle?

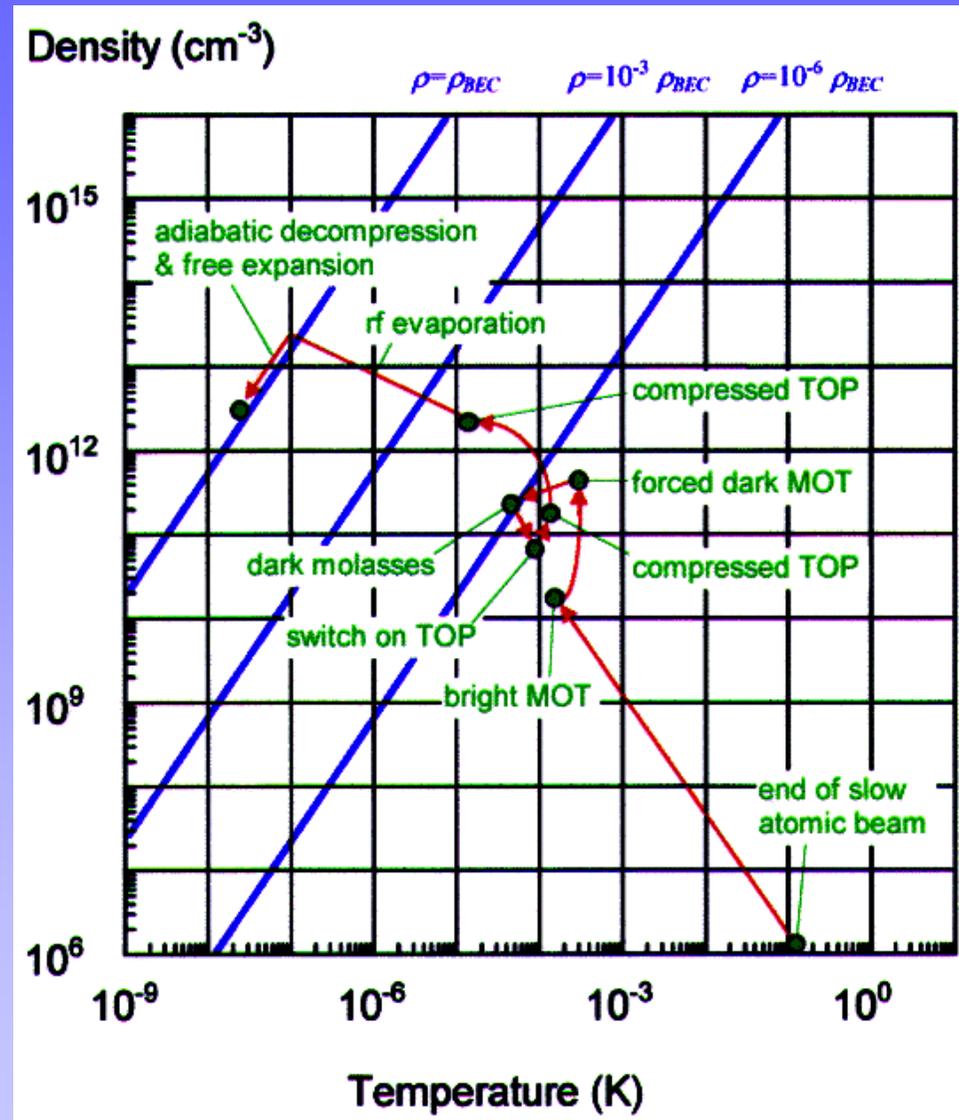
Thermal energy $k_B T$ smaller than the internal excitation of the particle

⇒ internal degrees of freedom are frozen out and do not matter for the thermodynamics

Interaction energy smaller than the internal excitation energy

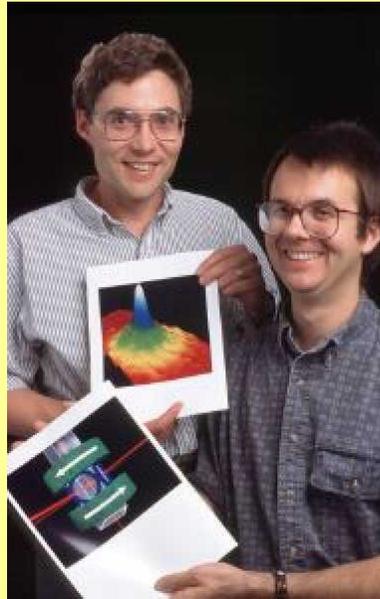
⇒ collisions do not remove or excite the bound electrons

The long, long road to BEC



BEC in Boulder, Juni 1995
(Rubidium)

Bose-Einstein condensation



Carl Wieman,
Eric Cornell



Wolfgang
Ketterle

Physics Nobel Prize 2001

1.0 mm

$T < T_c$

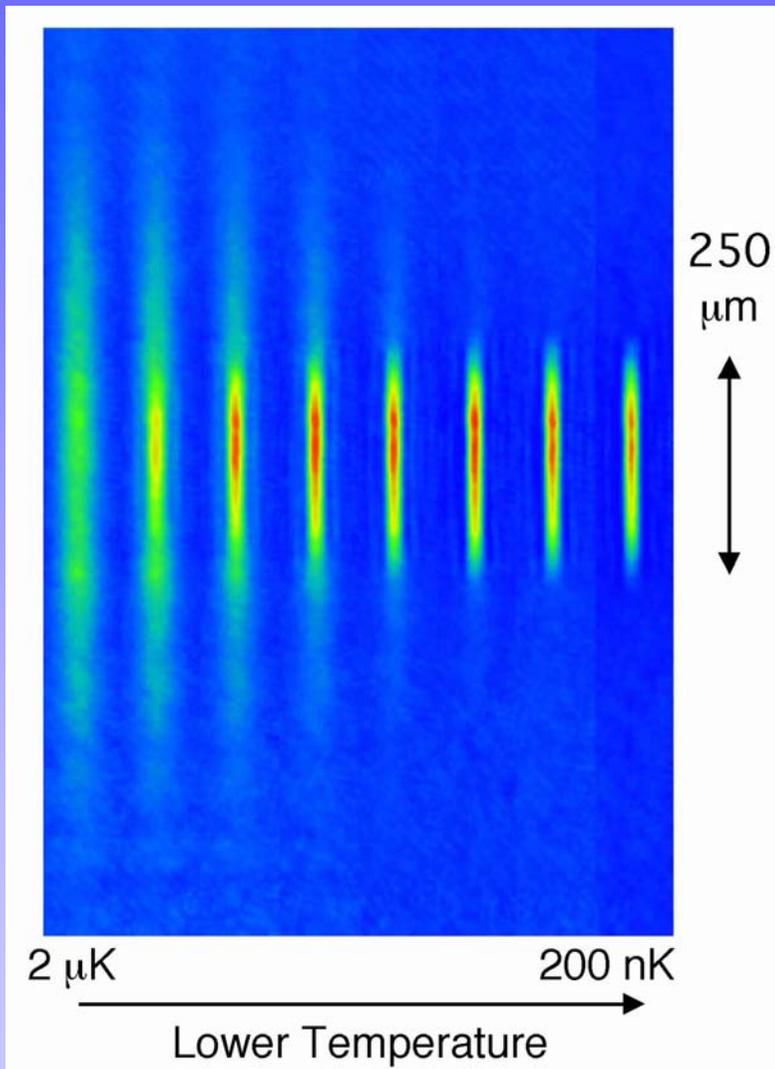
BEC at MIT, Nov. 1995 (Natrium)

BEC - “molecule” of the year

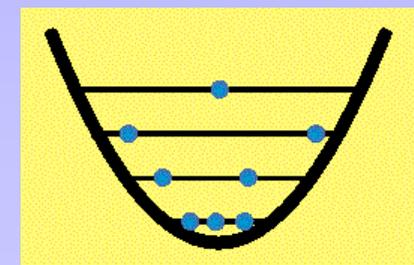
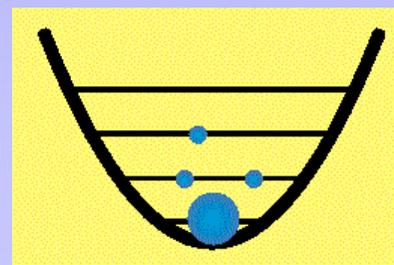
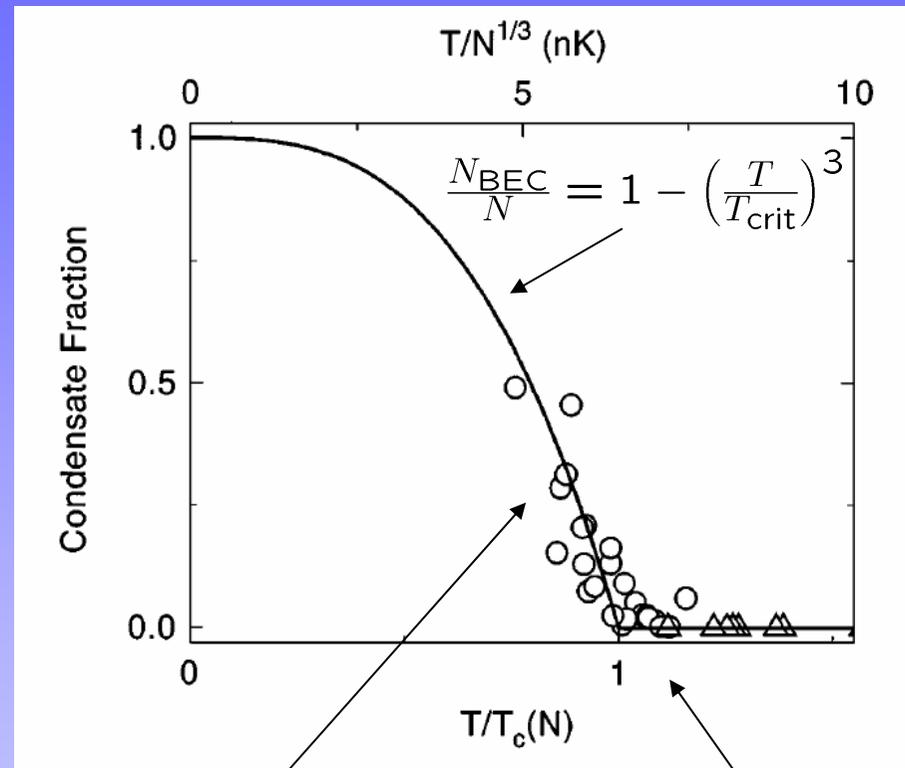


Transition to BEC

In-situ measurement



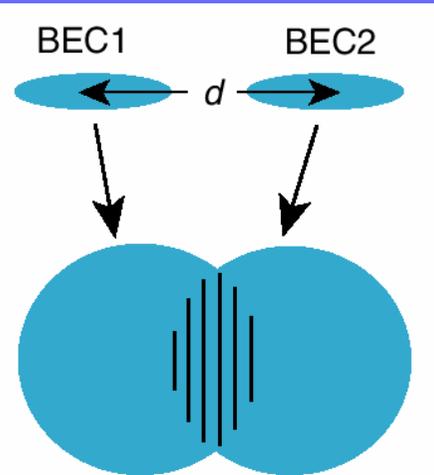
Number of atoms in the condensate



W. Ketterle *et al.* @ MIT

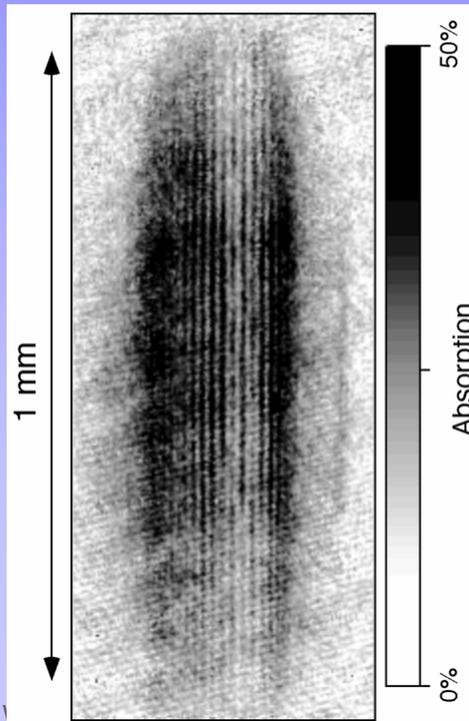
Matter-wave interference

trapped BECs

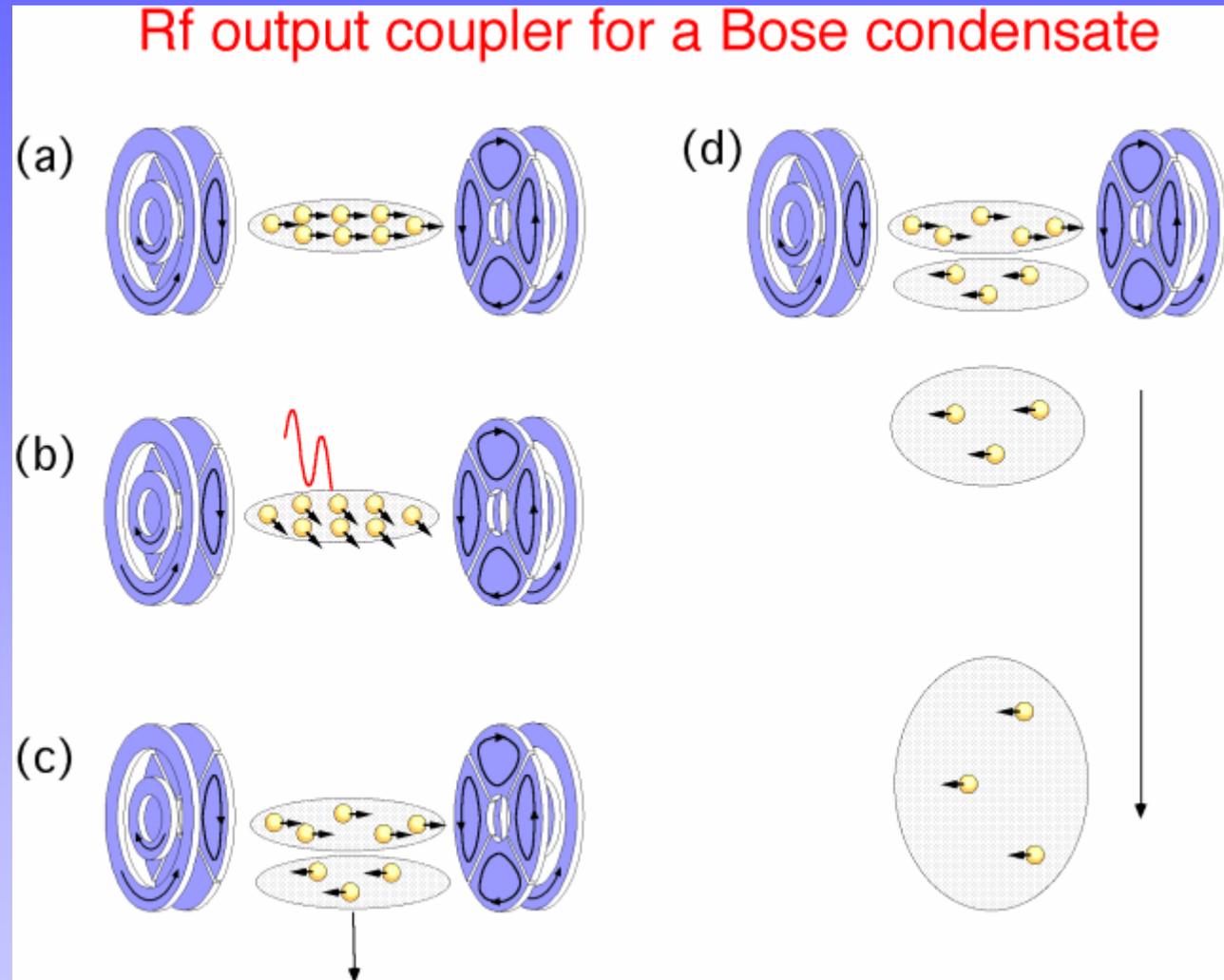
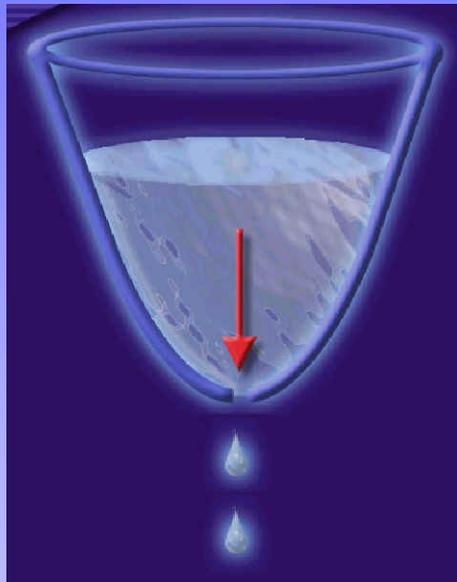


BECs after expansion
(expansion time Δt)

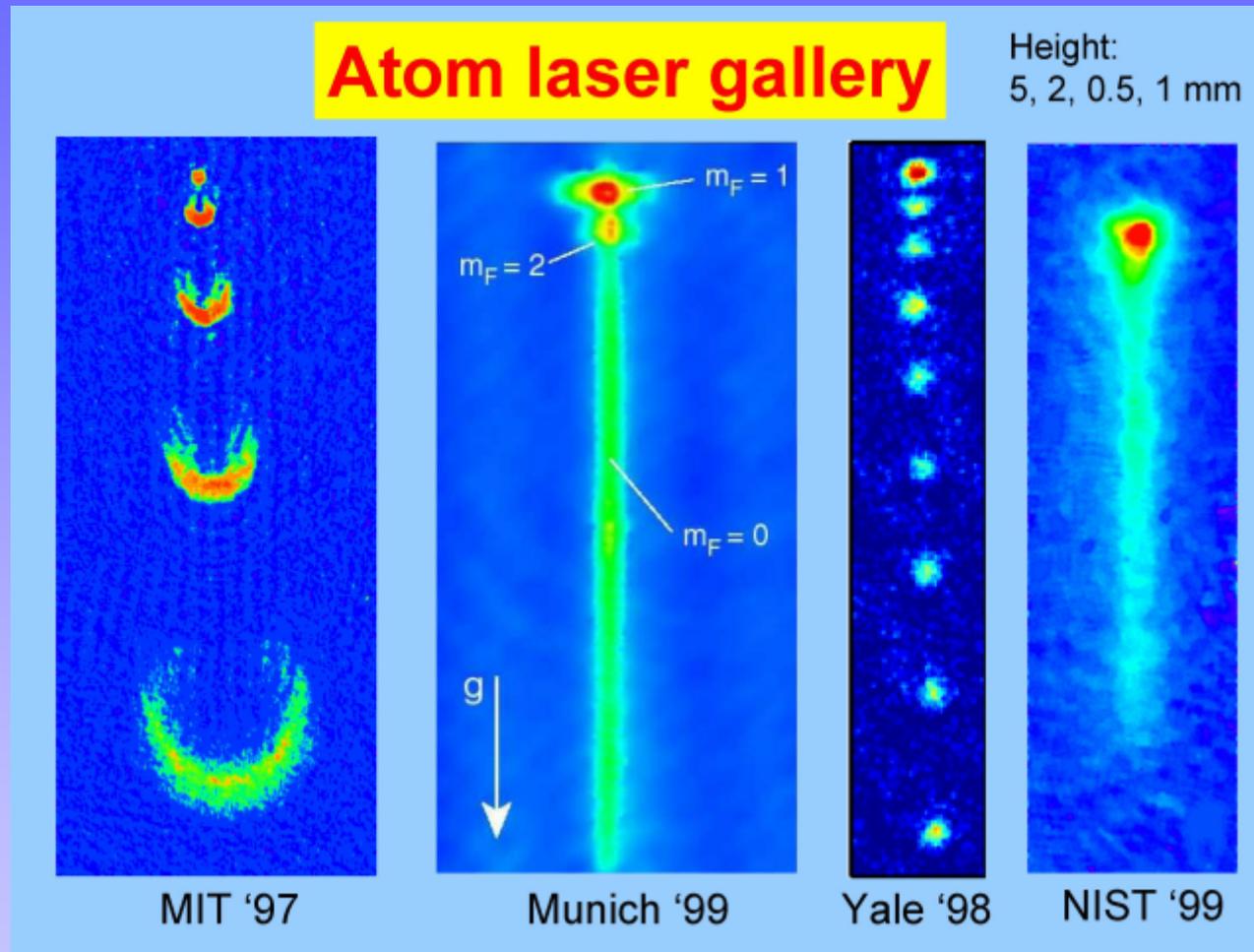
$$\Lambda_{dB} = \frac{h}{m \Delta v} = \frac{h \Delta t}{m d}$$



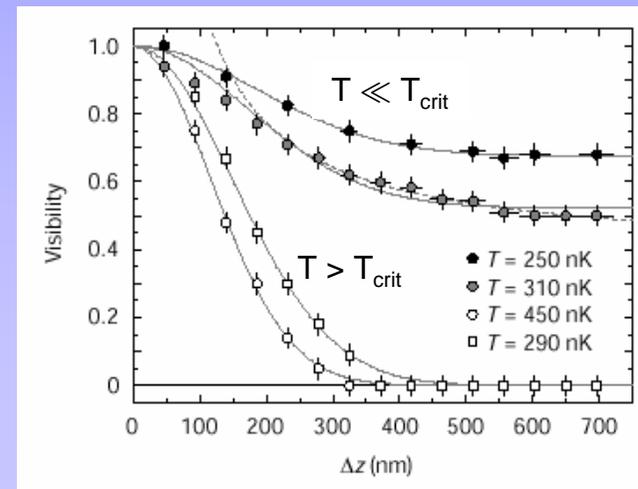
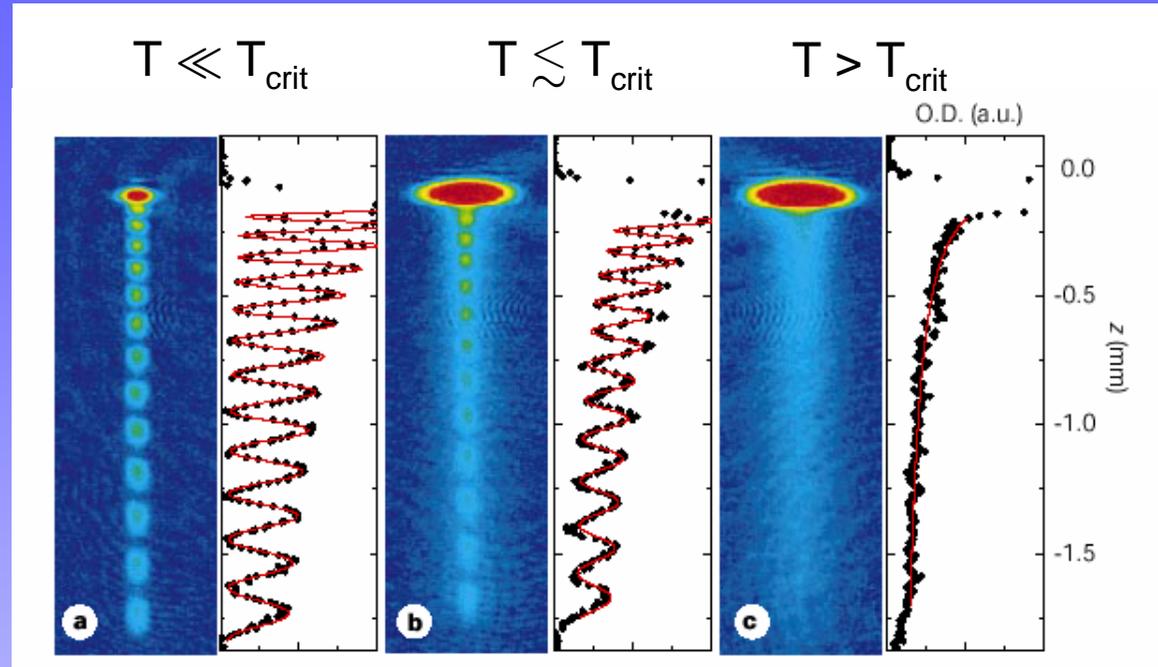
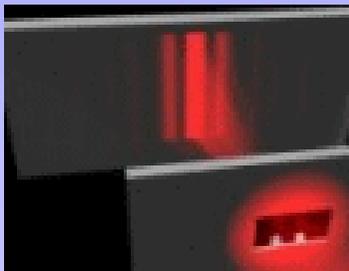
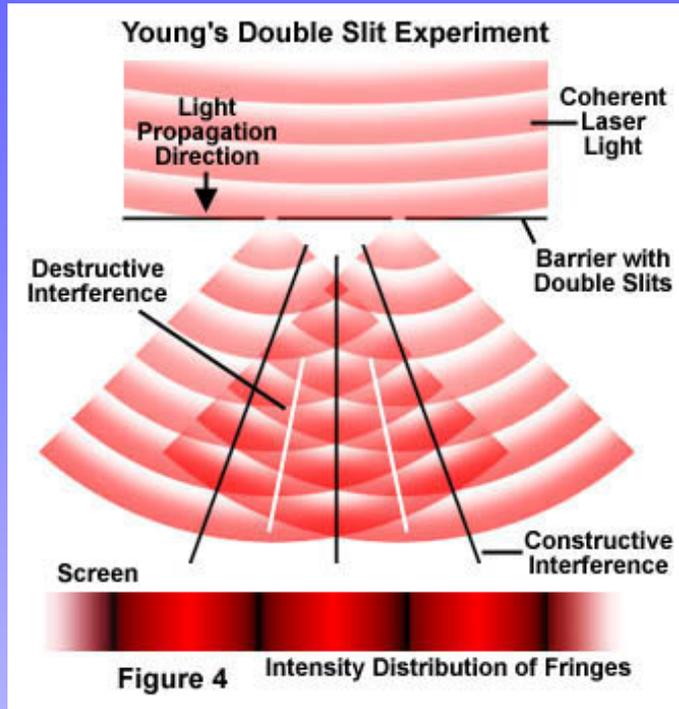
Output Coupler for a Bose condensate



Atom Lasers



Coherence of the atom laser



Gross-Pitaevskii equation

Gross-Pitaevskii equation for BEC mean-field interaction

$$i\hbar \frac{d\psi}{dt} = -\frac{\hbar^2}{2m} \nabla^2 \psi + U(r)\psi + g |\psi|^2 \psi$$

$A < 0$ attraction

$A = 0$ ideal gas

$A > 0$ repulsion

$$g = \frac{4\pi \hbar^2 A}{m}$$

ensemble creates „mean field“ (prop. to number density), which gives additional „potential“ in Schrödinger equation

→ simple „mean-field theory“ with nonlinear behavior
many phenomena (expansion, sound, collective oscillations) can be understood in this way !!

Thomas-Fermi regime

stationary solution

large ensemble

Gross-Pitaevskii equation for BEC mean-field interaction

$$i\hbar \frac{d\psi}{dt} = -\frac{\hbar^2}{2m} \nabla^2 \psi + U(r)\psi + g |\psi|^2 \psi$$

$A < 0$ attraction

$A = 0$ ideal gas

$A > 0$ repulsion

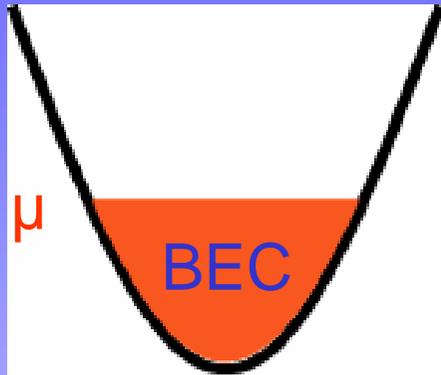
$$g = \frac{4\pi \hbar^2 A}{m}$$

very simple solution

$$n(\vec{r}) \propto \mu - U(\vec{r})$$

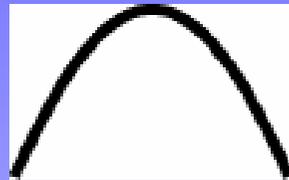
Thomas-Fermi limit

$U(r)$



harmonic
potential

$$n(r) \sim \mu - U(r)$$



mean-field
potential

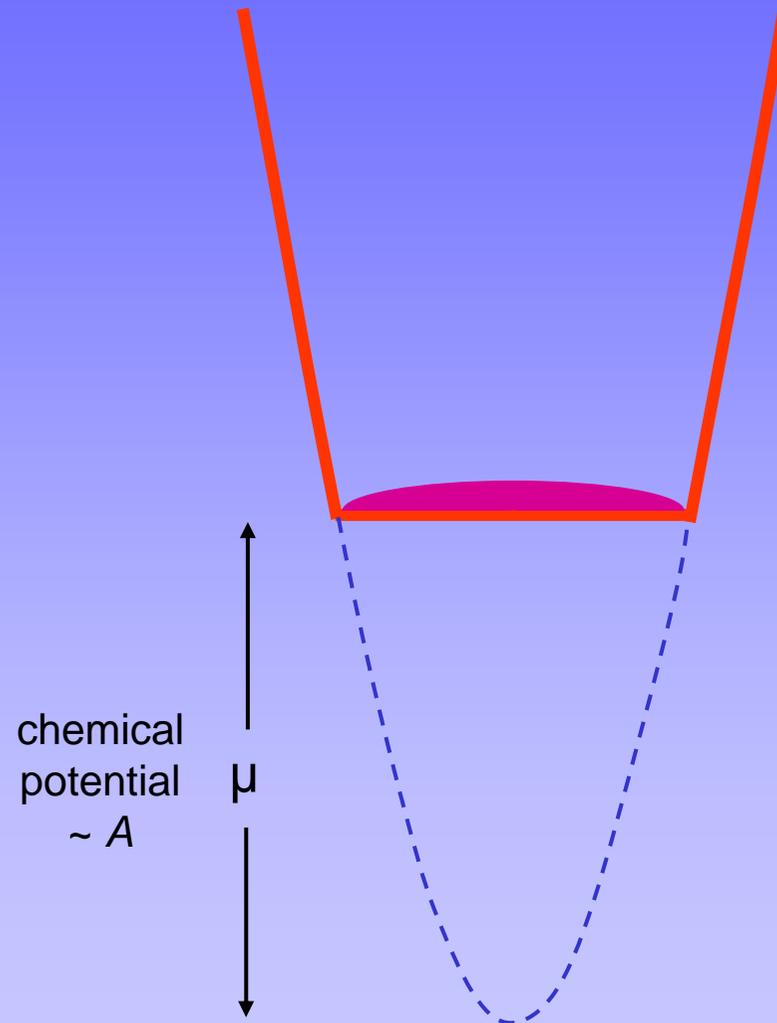
=



total potential

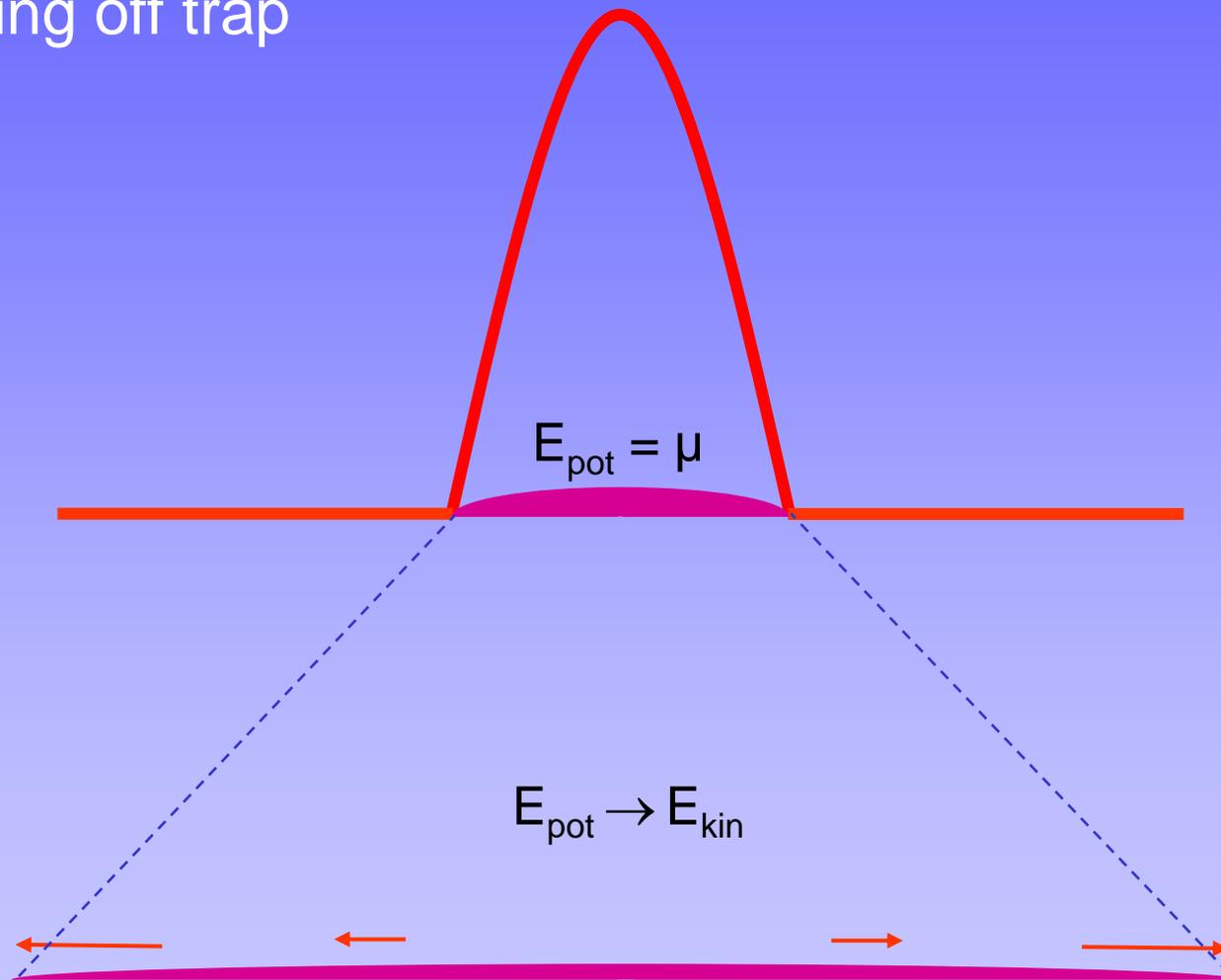
BEC density distribution is inverse shape of trap potential

Effect of interactions



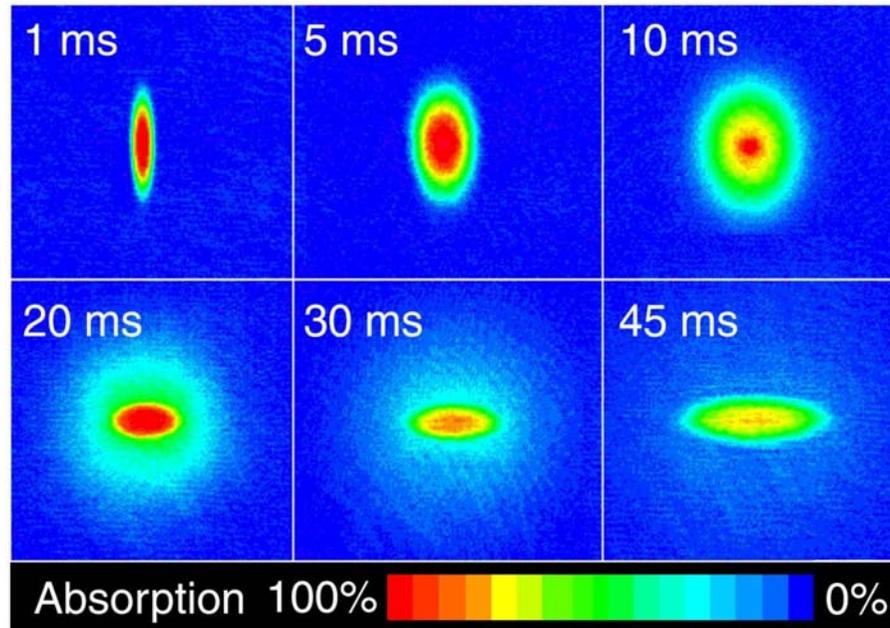
Expansion

switching off trap



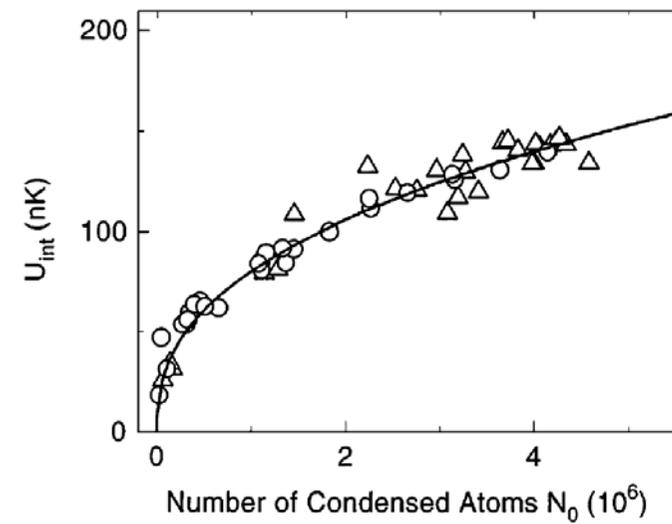
Influence of interactions

Free expansion of a Bose condensate



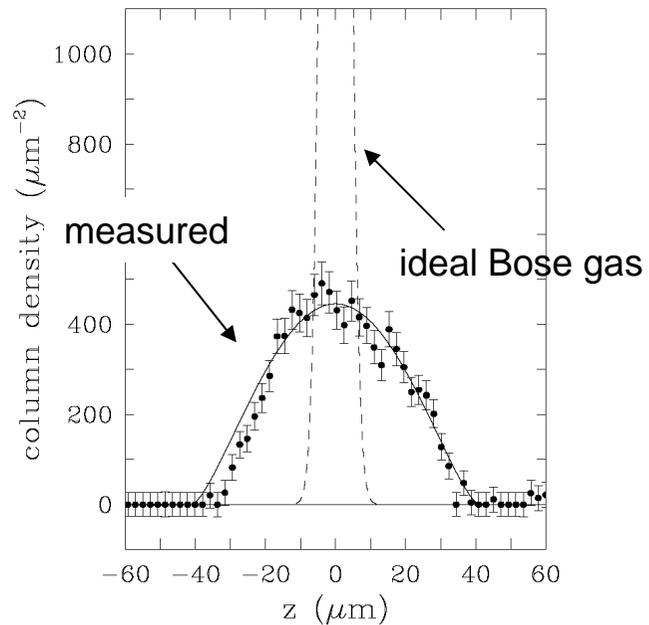
W. Ketterle *et al.* @ MIT

Internal energy



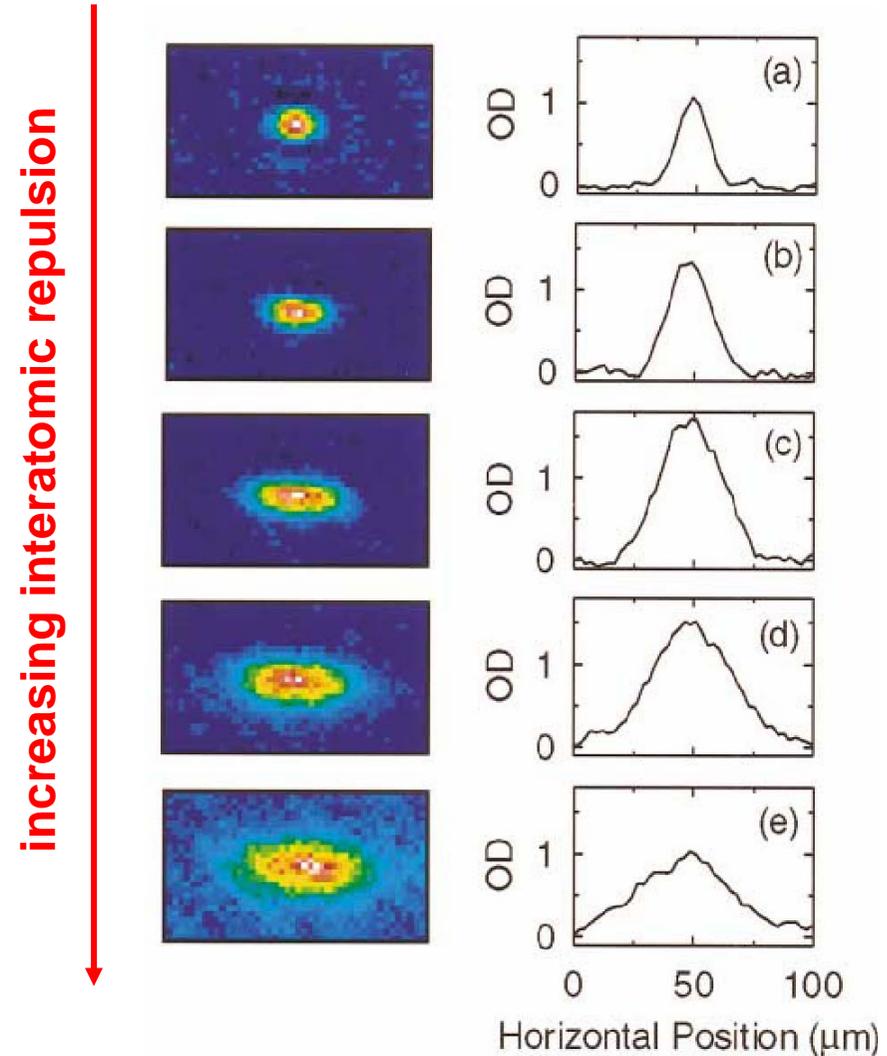
Influence of interactions (cont'd)

Density distribution



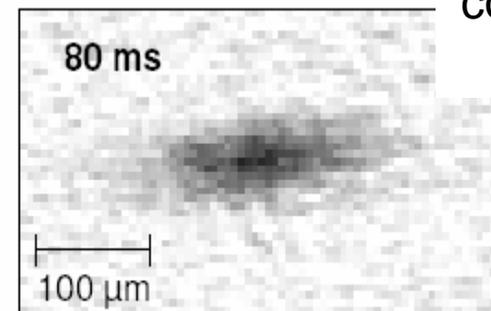
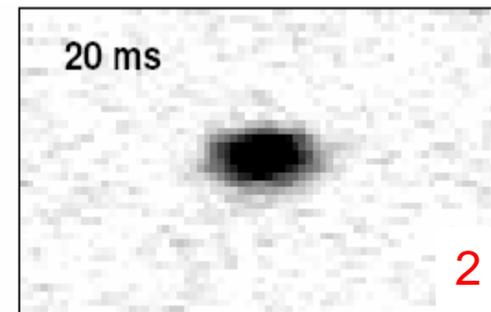
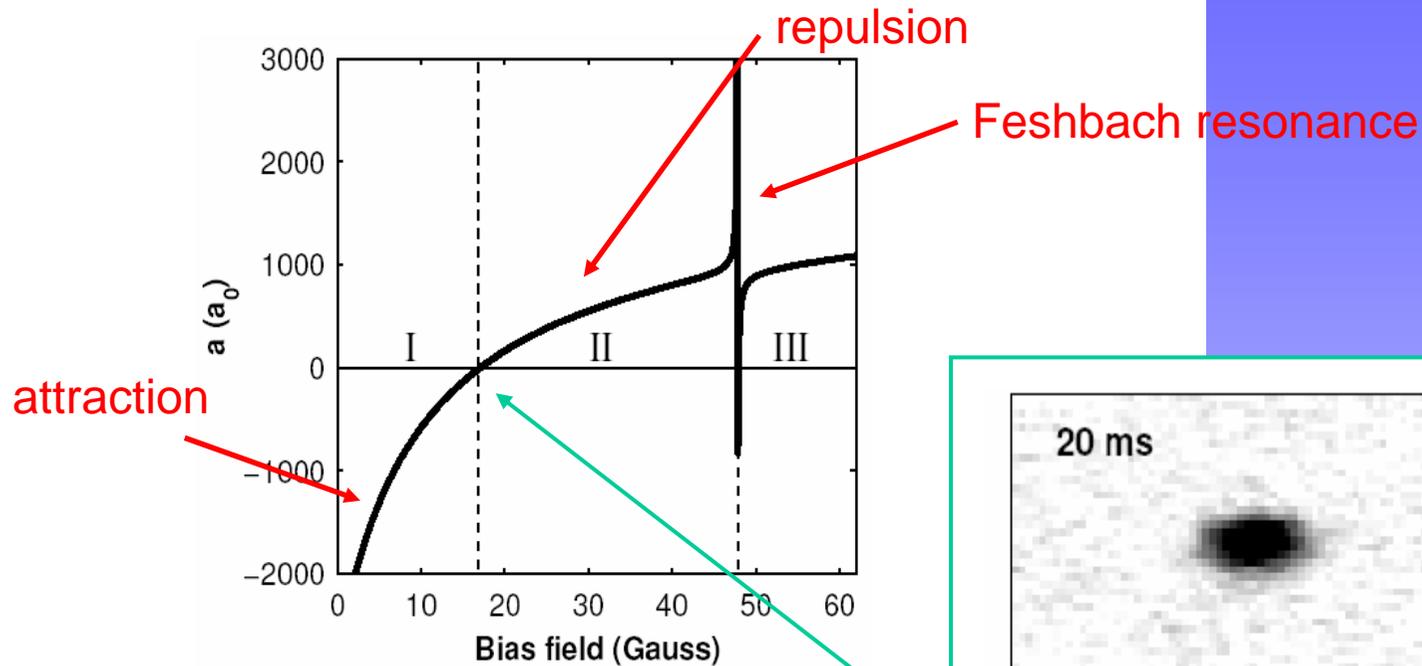
L. Hau *et al.* @ Harvard

Manipulation of the interaction strength



E. Cornell, C. Wieman *et al.* @ JILA Boulder

Changing the mean-field interaction

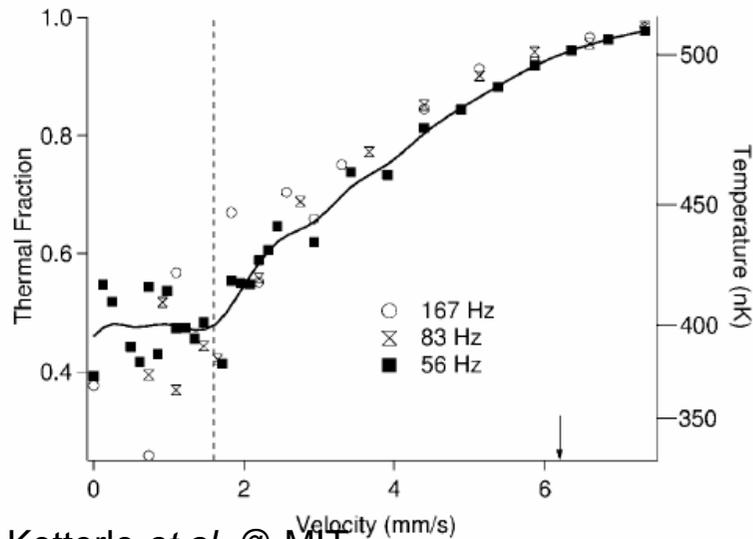
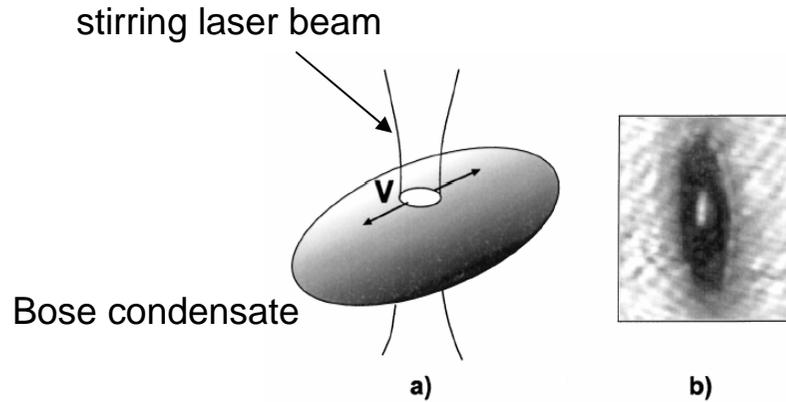


2 μm in 100 ms
= 7 cm/h
corresponds to
50 pK !!!

courtesy Rudi Grimm (Universität Innsbruck)

Superfluidity

Critical velocity



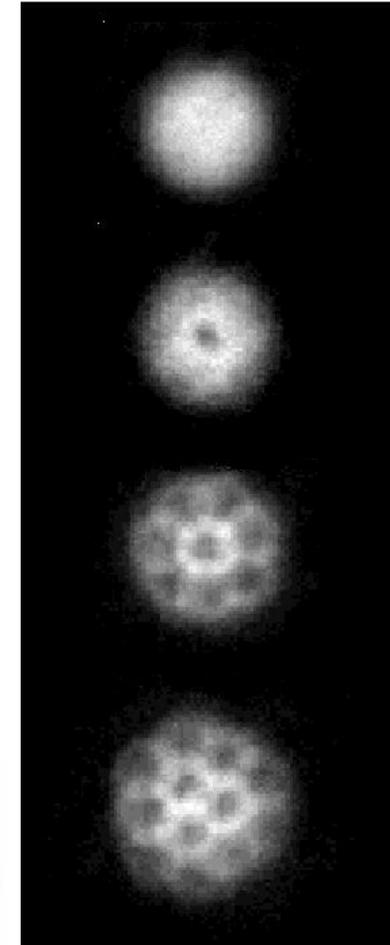
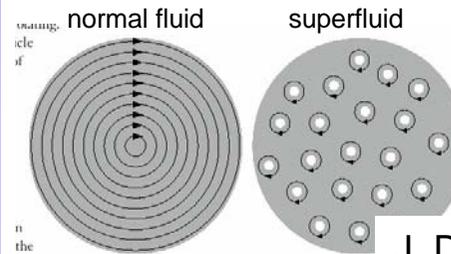
W. Ketterle *et al.* @ MIT

Quantized flux vortices



quantization of the velocity field due to macroscopic wavefunction

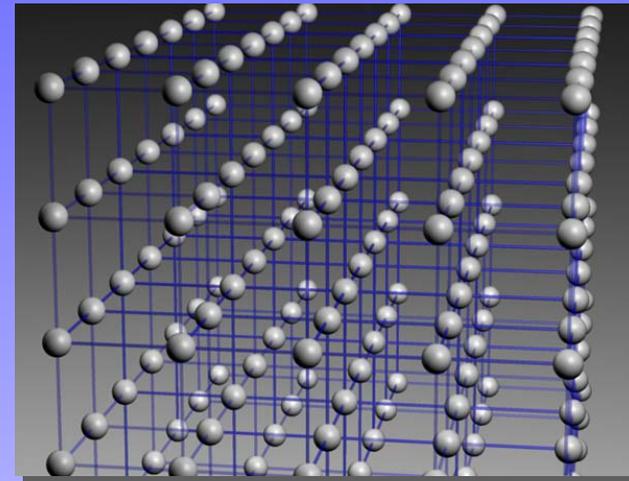
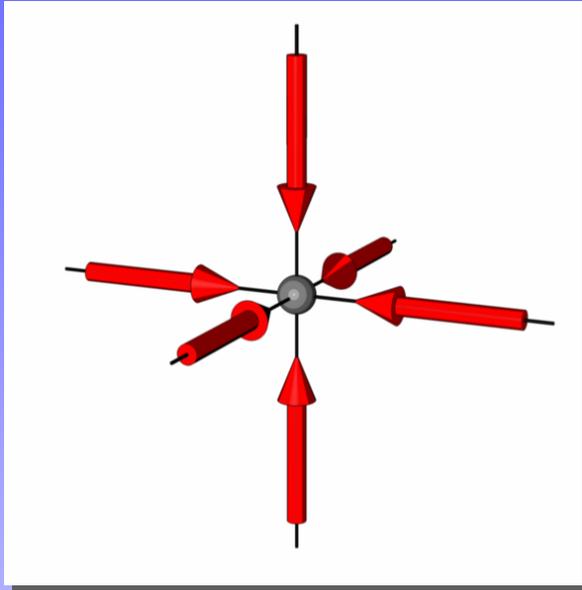
$$|\mathbf{v}| = n \frac{h}{m}$$



J. Dalibard *et al.* @ ENS, Paris

BEC in 3D Lattice Potential

courtesy Immanuel Bloch (Universität Mainz)



- **Resulting potential consists of a simple cubic lattice**
- **BEC coherently populates more than 100,000 lattice sites**

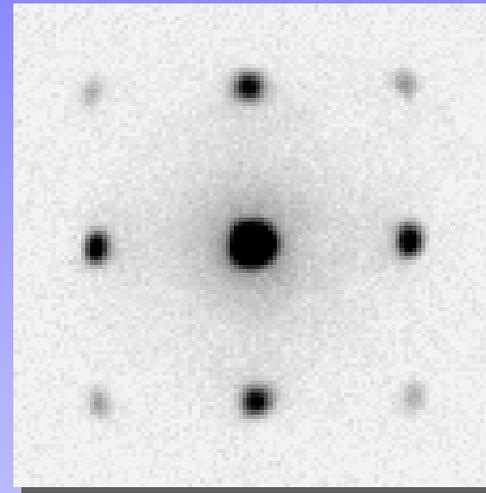
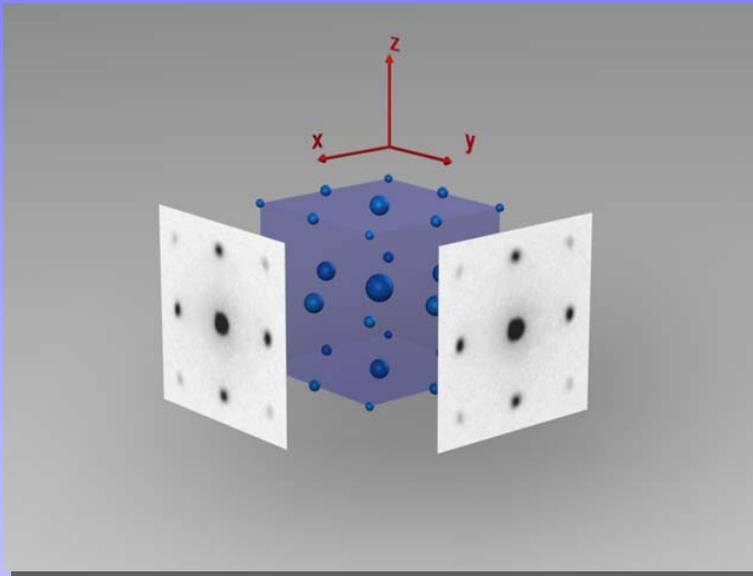
V_0 up to $22 E_{\text{recoil}}$

ω_r up to $2\pi \times 30 \text{ kHz}$

$n \approx 1-5$ atoms on average per site

Interference Pattern of matter waves

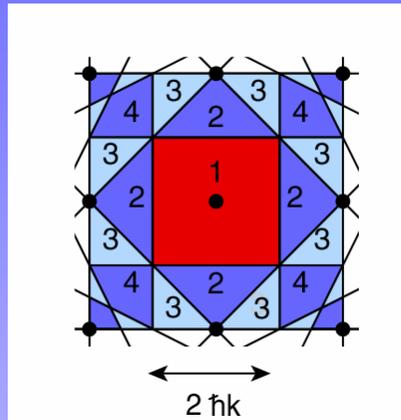
courtesy Immanuel Bloch (Universität Mainz)



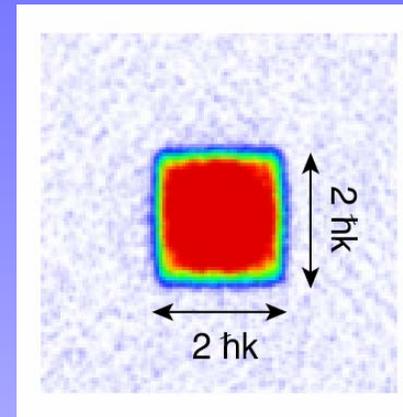
Mapping Brillouin zones

courtesy Immanuel Bloch (Universität Mainz)

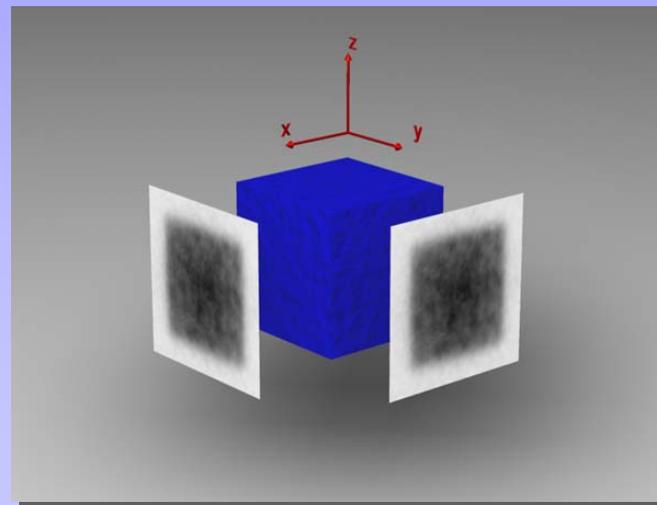
Brillouin Zones in 2D



Momentum distribution of a dephased condensate after turning off the lattice potential adiabatically



2D



3D

M. Greiner *et al.*, PRL **87**, 160405 (2001)

Basic idea of a Mott insulator

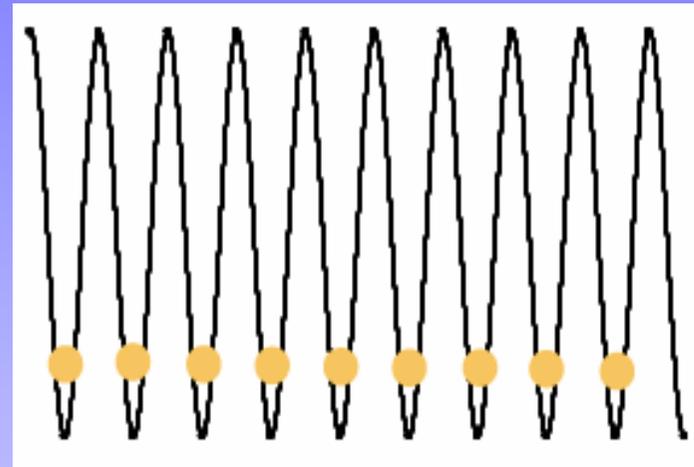
two different quantum phases ($T=0$)
separated by a *quantum phase transition*

BEC (superfluid)



- strong tunnel coupling
- fixed phase relations
- fluctuations of site occupation numbers

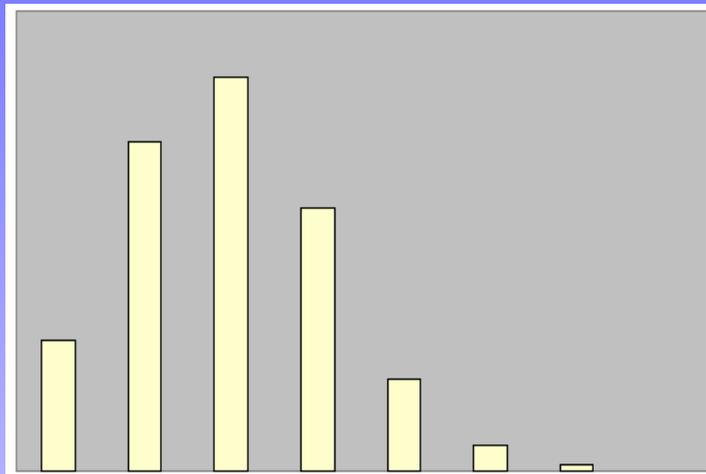
Mott insulator



- weak tunnel coupling
- no phase relations
- no fluctuations of site occupation numbers

Number fluctuation per lattice site

superfluid



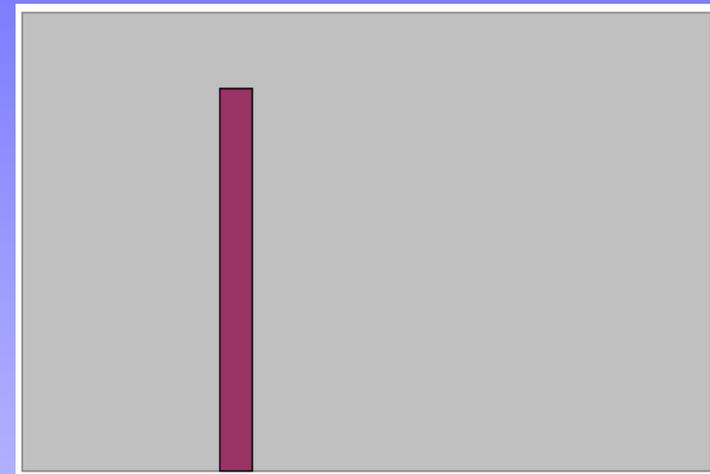
0 1 2 3 4 5 6

coherent state

(interference possible)

$$\langle n \rangle = 2$$

Mott insulator

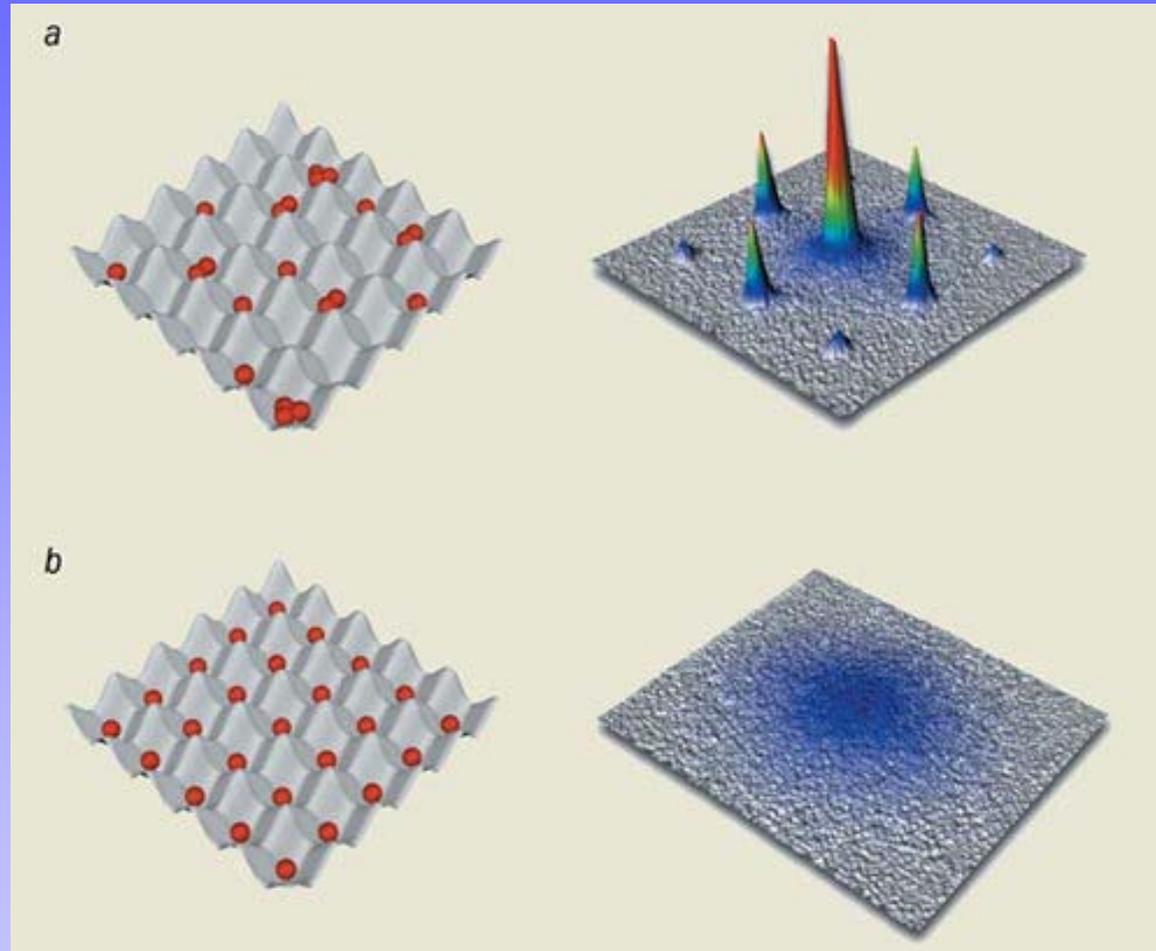


0 1 2 3 4 5 6

number state

(no interference)

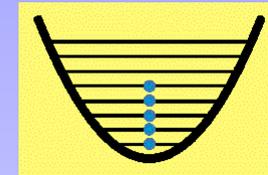
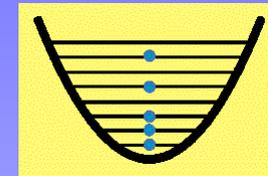
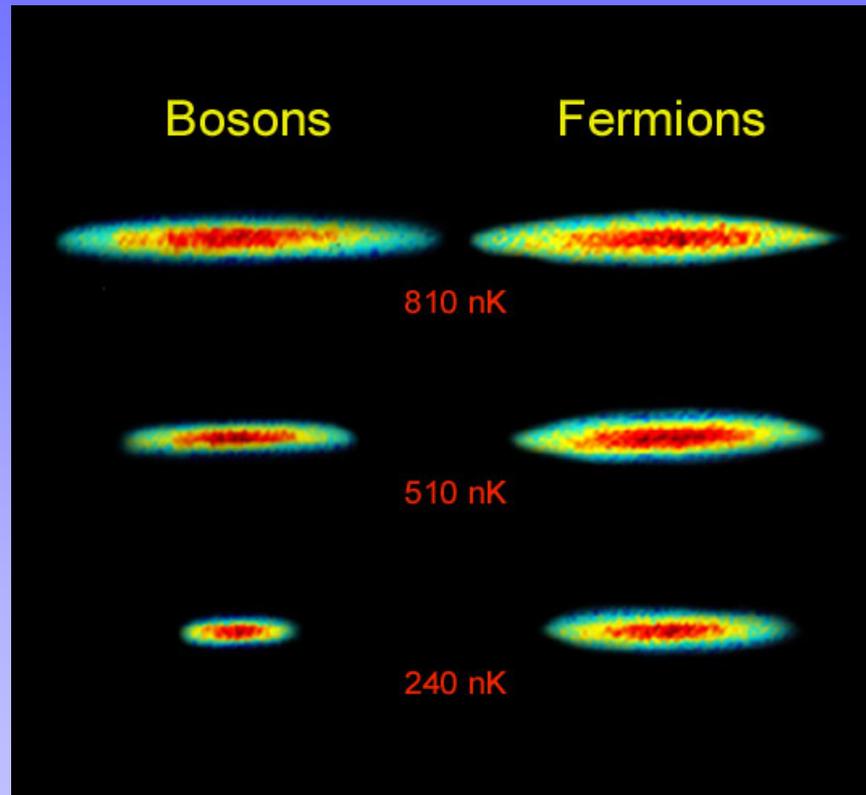
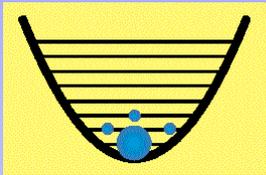
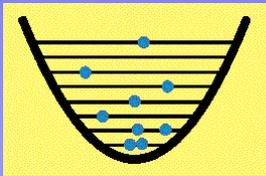
Observation of the Mott insulator



M. Greiner et al., Nature **415**, 39 (2002)
I. Bloch, Physics World, April 2004

Degenerate Fermi gases

Atomic white dwarf (Pauli pressure)



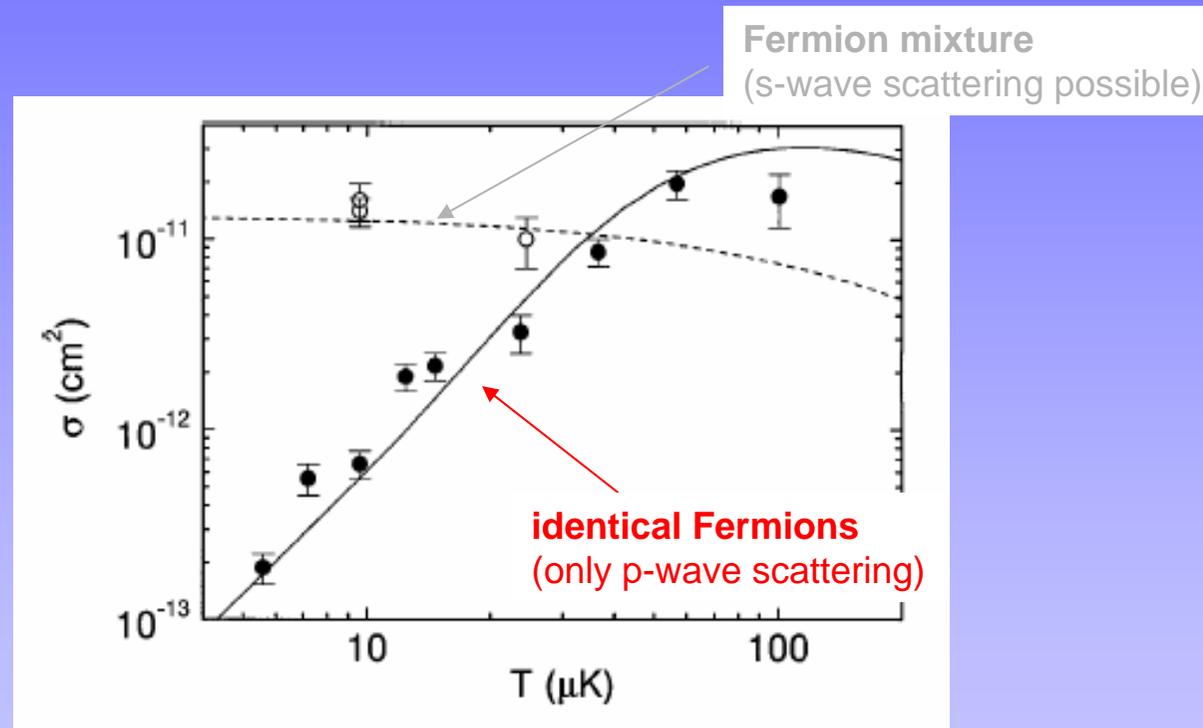
R. Hulet *et al.* © Rice University, Houston

Fermions

Suppression of elastic collisions

identical Fermions:

s-wave scattering length $A = 0$

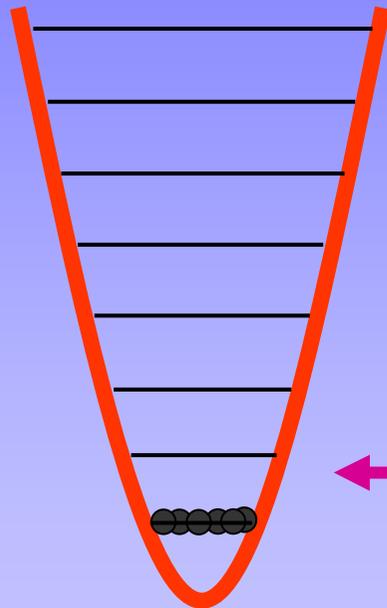


B. deMarco *et al.*, PRL **82**, 4208 (1999)

Two classes

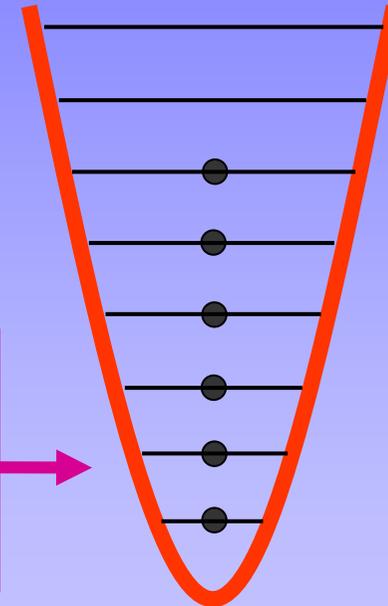
courtesy Rudi Grimm (Universität Innsbruck)

Bosons
integer spin



all in ground state:
Bose-Einstein condensate

Fermions
half-integer spin



only one particle per state:
degenerate Fermi gas

trapped atoms
at $T=0$

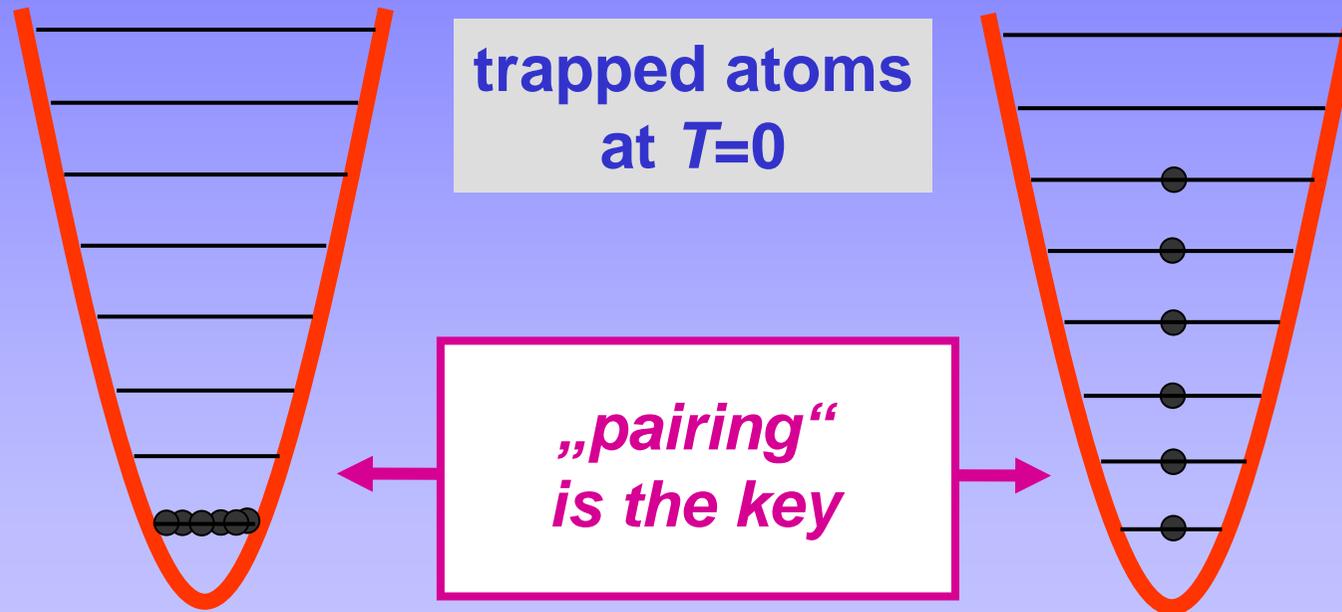
*these two
worlds
are connected !*

Two classes

courtesy Rudi Grimm (Universität Innsbruck)

Bosons
integer spin

Fermions
half-integer spin



all in ground state:
Bose-Einstein condensate

only one particle per state:
degenerate Fermi gas

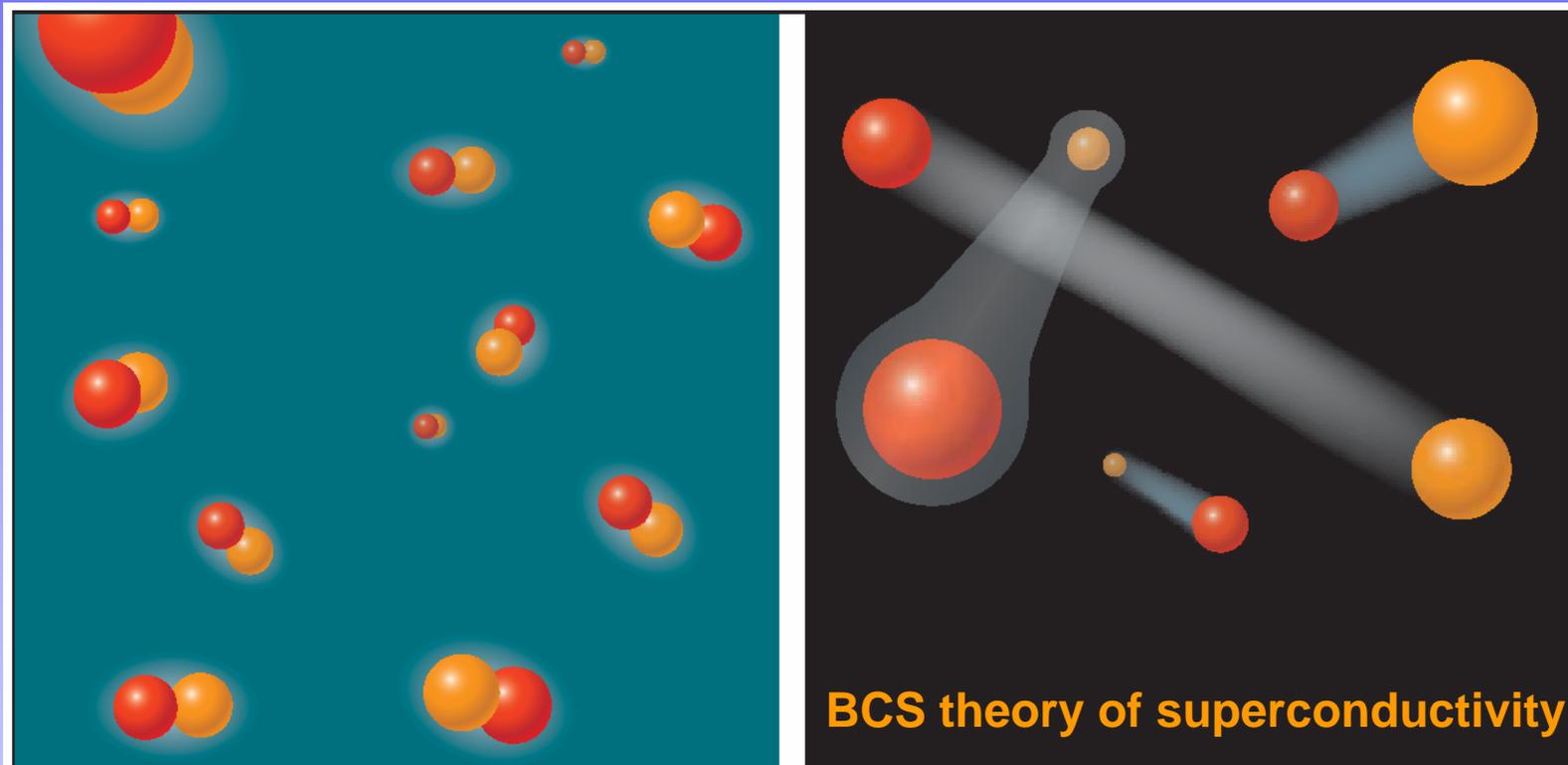
BEC-BCS crossover

courtesy Rudi Grimm (Universität Innsbruck)

molecules

crossover

Cooper pairs



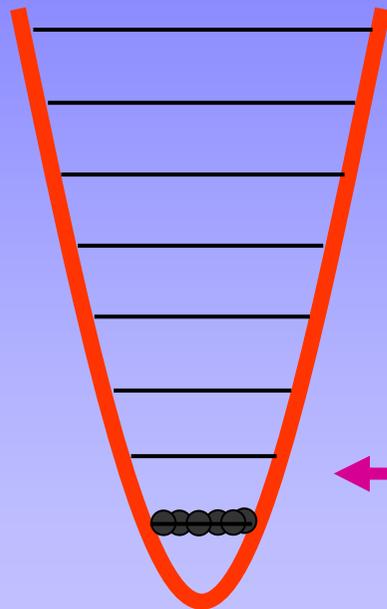
Tango or twist? In a magnetic field, atoms in different spin states can form molecules (*left*). Vary the field, and they might also form loose-knit Cooper pairs.

A. Cho, *Science* **301**, 750 (2003)

Two classes

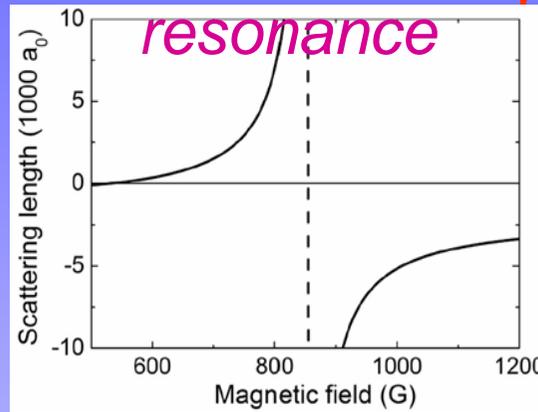
courtesy Rudi Grimm (Universität Innsbruck)

Bosons
integer spin

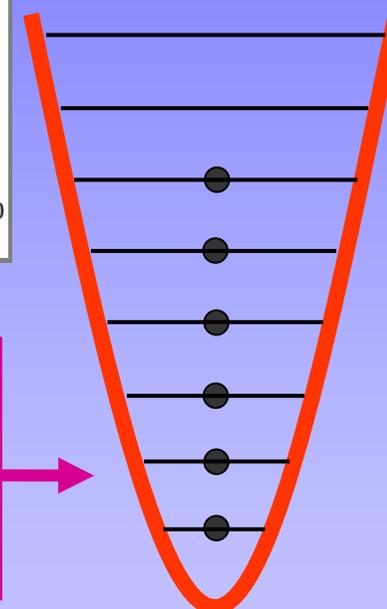


all in ground state:
Bose-Einstein condensate

Feshbach resonance



Fermions
half-integer spin



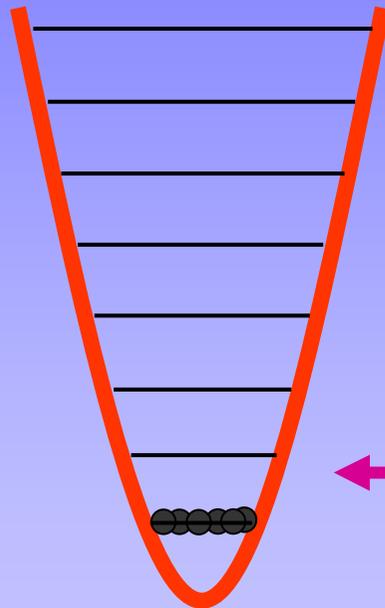
only one particle per state:
degenerate Fermi gas

interaction control !!!

Two classes

courtesy Rudi Grimm (Universität Innsbruck)

Bosons
integer spin



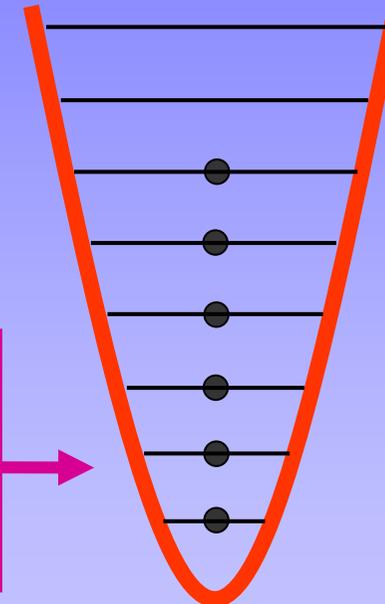
all in ground state:
Bose-Einstein condensate

neutron star



**exotic states
of matter**

Fermions
half-integer spin



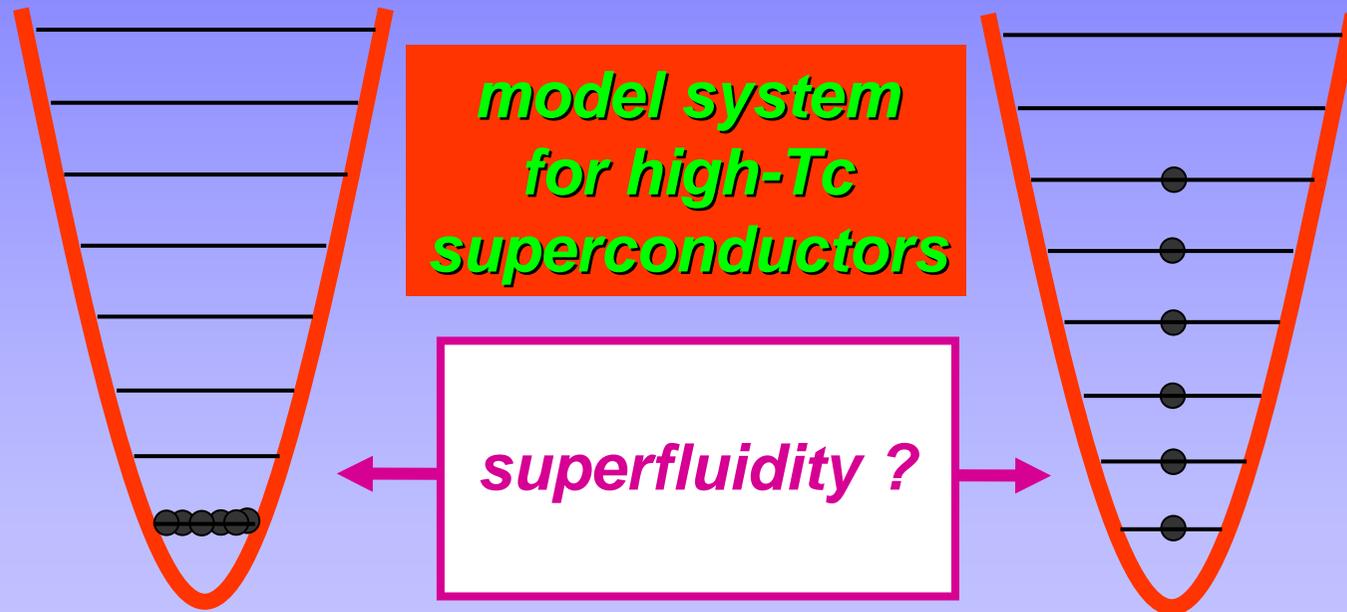
only one particle per state:
degenerate Fermi gas

Two classes

courtesy Rudi Grimm (Universität Innsbruck)

Bosons
integer spin

Fermions
half-integer spin

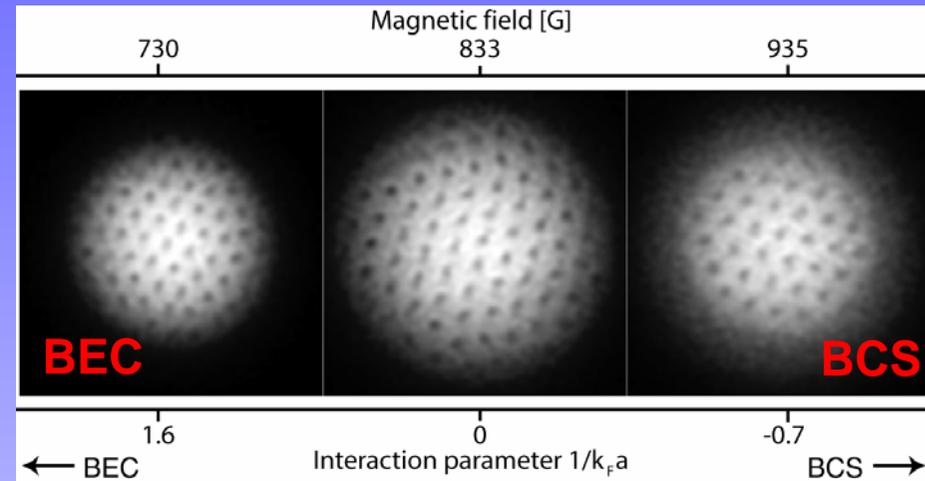
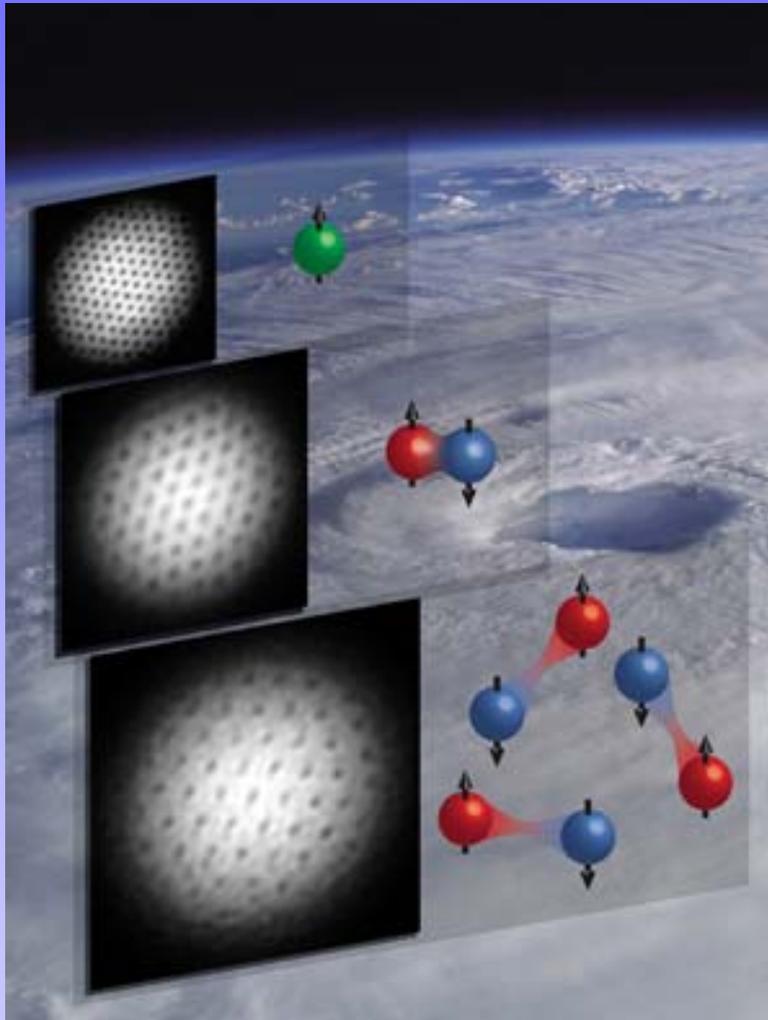


all in ground state:
Bose-Einstein condensate

only one particle per state:
degenerate Fermi gas

Superfluidity in a paired Fermi gas

Vortices in the BEC and BCS phase



M.W. Zwierlein *et al.*, Nature **435**, 1047 (2005)

Summary of Lecture 2

➤ **Quantum statistics of Bosons and Fermions**

- quantum statistics of composite particles
- Bose-Einstein condensation and Fermi degeneracy

➤ **Bose-Einstein condensation**

- giant matter-wave interference
- role of elastic collisions (mean-field interaction, superfluidity)
- Bose condensates in optical lattices (diffraction, Mott insulator)

➤ **Degenerate Fermi gases**

- suppression of s-wave scattering
- Pauli pressure
- BEC-BCS crossover